

Index theory and boundary value problems for general first-order elliptic differential operators

Lashi Bandara
(joint with Christian Bär)

Department of Mathematics
Brunel University London

10 May 2022

Motivation from index theory

D be a differential operator, seen as an unbounded operator on L^2 .

Motivation from index theory

D be a differential operator, seen as an unbounded operator on L^2 .

D is *Fredholm*: $\text{ran } D$ closed and $\text{ker } D$ and $\text{ker } D^*$ are finite dimensional.

Motivation from index theory

D be a differential operator, seen as an unbounded operator on L^2 .

D is *Fredholm*: $\text{ran } D$ closed and $\text{ker } D$ and $\text{ker } D^*$ are finite dimensional.

Index (analytic) of D :

$$\text{index } D = \dim \text{ker } D - \dim \text{ker } D^*.$$

Motivation from index theory

D be a differential operator, seen as an unbounded operator on L^2 .

D is *Fredholm*: $\text{ran } D$ closed and $\text{ker } D$ and $\text{ker } D^*$ are finite dimensional.

Index (analytic) of D :

$$\text{index } D = \dim \text{ker } D - \dim \text{ker } D^*.$$

Invariant, in particular, $\text{index } D_t = \text{index } D$ for continuous deformation of t through Fredholm operators.

Motivation from index theory

D be a differential operator, seen as an unbounded operator on L^2 .

D is *Fredholm*: $\text{ran } D$ closed and $\ker D$ and $\ker D^*$ are finite dimensional.

Index (analytic) of D :

$$\text{index } D = \dim \ker D - \dim \ker D^*.$$

Invariant, in particular, $\text{index } D_t = \text{index } D$ for continuous deformation of t through Fredholm operators.

Usually, D determined by geometry

Motivation from index theory

D be a differential operator, seen as an unbounded operator on L^2 .

D is *Fredholm*: $\text{ran } D$ closed and $\ker D$ and $\ker D^*$ are finite dimensional.

Index (analytic) of D :

$$\text{index } D = \dim \ker D - \dim \ker D^*.$$

Invariant, in particular, $\text{index } D_t = \text{index } D$ for continuous deformation of t through Fredholm operators.

Usually, D determined by geometry : $t \rightarrow D_t$ geometric deformation

Motivation from index theory

D be a differential operator, seen as an unbounded operator on L^2 .

D is *Fredholm*: $\text{ran } D$ closed and $\ker D$ and $\ker D^*$ are finite dimensional.

Index (analytic) of D :

$$\text{index } D = \dim \ker D - \dim \ker D^*.$$

Invariant, in particular, $\text{index } D_t = \text{index } D$ for continuous deformation of t through Fredholm operators.

Usually, D determined by geometry : $t \rightarrow D_t$ geometric deformation, i.e., evolving time slices in spacetime.

Index formula: relate geometry, topology and boundary.

Index formula: relate geometry, topology and boundary.

$$\text{index } D = \int \text{“Curvatures related to } D\text{”} + \text{“Boundary contribution”}.$$

Atiyah-Patodi-Singer

Index formula: relate geometry, topology and boundary.

$$\text{index } D = \int \text{“Curvatures related to } D\text{”} + \text{“Boundary contribution”}.$$

Atiyah-Patodi-Singer (sufficiently nice D , in particular elliptic):

$$\text{“Boundary contribution”} = \frac{\ker(A) + \eta(A)}{2},$$

Index formula: relate geometry, topology and boundary.

$$\text{index } D = \int \text{“Curvatures related to } D\text{”} + \text{“Boundary contribution”}.$$

Atiyah-Patodi-Singer (sufficiently nice D , in particular elliptic):

$$\text{“Boundary contribution”} = \frac{\ker(A) + \eta(A)}{2},$$

A adapted operator to the boundary - determined by D

Index formula: relate geometry, topology and boundary.

$$\text{index } D = \int \text{“Curvatures related to } D\text{”} + \text{“Boundary contribution”}.$$

Atiyah-Patodi-Singer (sufficiently nice D , in particular elliptic):

$$\text{“Boundary contribution”} = \frac{\ker(A) + \eta(A)}{2},$$

A adapted operator to the boundary - determined by D ,

$\eta(A)$ - measuring spectral *asymmetry* of A .

Index formula: relate geometry, topology and boundary.

$$\text{index } D = \int \text{“Curvatures related to } D\text{”} + \text{“Boundary contribution”}.$$

Atiyah-Patodi-Singer (sufficiently nice D , in particular elliptic):

$$\text{“Boundary contribution”} = \frac{\ker(A) + \eta(A)}{2},$$

A adapted operator to the boundary - determined by D ,

$\eta(A)$ - measuring spectral *asymmetry* of A .

Need: boundary condition to make formula work.

Index formula: relate geometry, topology and boundary.

$$\text{index } D = \int \text{“Curvatures related to } D\text{”} + \text{“Boundary contribution”}.$$

Atiyah-Patodi-Singer (sufficiently nice D , in particular elliptic):

$$\text{“Boundary contribution”} = \frac{\ker(A) + \eta(A)}{2},$$

A adapted operator to the boundary - determined by D ,

$\eta(A)$ - measuring spectral *asymmetry* of A .

Need: boundary condition to make formula work.

Non-local boundary conditions

Index formula: relate geometry, topology and boundary.

$$\text{index } D = \int \text{“Curvatures related to } D\text{”} + \text{“Boundary contribution”}.$$

Atiyah-Patodi-Singer (sufficiently nice D , in particular elliptic):

$$\text{“Boundary contribution”} = \frac{\ker(A) + \eta(A)}{2},$$

A adapted operator to the boundary - determined by D ,

$\eta(A)$ - measuring spectral *asymmetry* of A .

Need: boundary condition to make formula work.

Non-local boundary conditions: topologically obstructed for local boundary conditions.

Motivation: a Euclidean example

$\Omega \subset \mathbb{R}^3$, domain with smooth compact boundary $\partial\Omega \subset \Omega$.

Motivation: a Euclidean example

$\Omega \subset \mathbb{R}^3$, domain with smooth compact boundary $\partial\Omega \subset \Omega$.

Spin “bundle” $\Delta\Omega \cong \mathbb{C}^2$.

Motivation: a Euclidean example

$\Omega \subset \mathbb{R}^3$, domain with smooth compact boundary $\partial\Omega \subset \Omega$.

Spin “bundle” $\Delta\Omega \cong \mathbb{C}^2$.

Representation $\rho(e_i)u := \sigma_i u$, given by Pauli matrices:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Motivation: a Euclidean example

$\Omega \subset \mathbb{R}^3$, domain with smooth compact boundary $\partial\Omega \subset \Omega$.

Spin “bundle” $\Delta\Omega \cong \mathbb{C}^2$.

Representation $\rho(e_i)u := \sigma_i u$, given by Pauli matrices:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Spin-Dirac operator:

$$\not{D}f := -i \sum_{j=1}^3 \rho(e_j) \partial_j f = -i (\sigma_1 \partial_1 f + \sigma_2 \partial_2 f + \sigma_3 \partial_3 f).$$

Motivation: a Euclidean example

$\Omega \subset \mathbb{R}^3$, domain with smooth compact boundary $\partial\Omega \subset \Omega$.

Spin “bundle” $\Delta\Omega \cong \mathbb{C}^2$.

Representation $\rho(e_i)u := \sigma_i u$, given by Pauli matrices:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Spin-Dirac operator:

$$\not{D}f := -i \sum_{j=1}^3 \rho(e_j) \partial_j f = -i (\sigma_1 \partial_1 f + \sigma_2 \partial_2 f + \sigma_3 \partial_3 f).$$

Satisfies:

$$\not{D}^2 f = \sum_{k=1}^3 \sum_{j=1}^2 \partial_k^2 f_j e_j = \Delta f.$$

Adapted operator

$$\mathcal{D}_c, \text{ with } \text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \left\{ u \in C_c^\infty(\Omega) : \text{spt } u \subset \mathring{\Omega} \right\}.$$

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \left\{ u \in C_c^\infty(\Omega) : \text{spt } u \subset \overset{\circ}{\Omega} \right\}$.

\mathcal{D}_c is symmetric.

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \{u \in C_c^\infty(\Omega) : \text{spt } u \subset \mathring{\Omega}\}$.

\mathcal{D}_c is symmetric. Define:

$$\mathcal{D}_{\max} := \mathcal{D}_c^*, \quad \mathcal{D}_{\min} := \overline{\mathcal{D}_c} = (\mathcal{D}_c^*)^*.$$

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \{u \in C_c^\infty(\Omega) : \text{spt } u \subset \mathring{\Omega}\}$.

\mathcal{D}_c is symmetric. Define:

$$\mathcal{D}_{\max} := \mathcal{D}_c^*, \quad \mathcal{D}_{\min} := \overline{\mathcal{D}_c} = (\mathcal{D}_c^*)^*.$$

Adapted operator on the boundary:

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \{u \in C_c^\infty(\Omega) : \text{spt } u \subset \mathring{\Omega}\}$.

\mathcal{D}_c is symmetric. Define:

$$\mathcal{D}_{\max} := \mathcal{D}_c^*, \quad \mathcal{D}_{\min} := \overline{\mathcal{D}_c} = (\mathcal{D}_c^*)^*.$$

Adapted operator on the boundary: fix ν outward pointing normal,

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \{u \in C_c^\infty(\Omega) : \text{spt } u \subset \mathring{\Omega}\}$.

\mathcal{D}_c is symmetric. Define:

$$\mathcal{D}_{\max} := \mathcal{D}_c^*, \quad \mathcal{D}_{\min} := \overline{\mathcal{D}_c} = (\mathcal{D}_c^*)^*.$$

Adapted operator on the boundary: fix ν outward pointing normal, $\{\tilde{e}_i\}_{i=1}^2$ orthonormal vectors at $x \in \partial\Omega$.

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \left\{ u \in C_c^\infty(\Omega) : \text{spt } u \subset \overset{\circ}{\Omega} \right\}$.

\mathcal{D}_c is symmetric. Define:

$$\mathcal{D}_{\max} := \mathcal{D}_c^*, \quad \mathcal{D}_{\min} := \overline{\mathcal{D}_c} = (\mathcal{D}_c^*)^*.$$

Adapted operator on the boundary: fix ν outward pointing normal, $\{\tilde{e}_i\}_{i=1}^2$ orthonormal vectors at $x \in \partial\Omega$. Then,

$$A_0 u(x) := \rho(\nu)^{-1} \sum_{i=1}^2 \left(\rho(\tilde{e}_i) \nabla_{\tilde{e}_i}^{\mathbb{C}^2} u \right) (x).$$

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \{u \in C_c^\infty(\Omega) : \text{spt } u \subset \mathring{\Omega}\}$.

\mathcal{D}_c is symmetric. Define:

$$\mathcal{D}_{\max} := \mathcal{D}_c^*, \quad \mathcal{D}_{\min} := \overline{\mathcal{D}_c} = (\mathcal{D}_c^*)^*.$$

Adapted operator on the boundary: fix ν outward pointing normal, $\{\tilde{e}_i\}_{i=1}^2$ orthonormal vectors at $x \in \partial\Omega$. Then,

$$A_0 u(x) := \rho(\nu)^{-1} \sum_{i=1}^2 \left(\rho(\tilde{e}_i) \nabla_{\tilde{e}_i}^{\mathbb{C}^2} u \right) (x).$$

Elliptic regularity: $\text{dom}(A_0) = H^1(\partial\Omega; \mathbb{C}^2)$.

Adapted operator

\mathcal{D}_c , with $\text{dom}(\mathcal{D}_c) = C_{cc}^\infty(\Omega) := \{u \in C_c^\infty(\Omega) : \text{spt } u \subset \overset{\circ}{\Omega}\}$.

\mathcal{D}_c is symmetric. Define:

$$\mathcal{D}_{\max} := \mathcal{D}_c^*, \quad \mathcal{D}_{\min} := \overline{\mathcal{D}_c} = (\mathcal{D}_c^*)^*.$$

Adapted operator on the boundary: fix ν outward pointing normal, $\{\tilde{e}_i\}_{i=1}^2$ orthonormal vectors at $x \in \partial\Omega$. Then,

$$A_0 u(x) := \rho(\nu)^{-1} \sum_{i=1}^2 \left(\rho(\tilde{e}_i) \nabla_{\tilde{e}_i}^{\mathbb{C}^2} u \right) (x).$$

Elliptic regularity: $\text{dom}(A_0) = H^1(\partial\Omega; \mathbb{C}^2)$.

There exists $B \in \mathcal{L}(L^2(\partial\Omega; \mathbb{C}^2))$ such that $A := A_0 + B$ self-adjoint.

Boundary conditions

Boundary conditions for \mathcal{D} “live” in:

Boundary conditions

Boundary conditions for \mathcal{D} “live” in:

$$\check{H}(A) := \chi_{(-\infty, 0]}(A)H^{\frac{1}{2}}(\partial\Omega; \mathbb{C}^2) \oplus \chi_{(0, \infty)}(A)H^{-\frac{1}{2}}(\partial\Omega; \mathbb{C}^2).$$

Boundary conditions

Boundary conditions for \mathcal{D} “live” in:

$$\check{H}(A) := \chi_{(-\infty, 0]}(A)H^{\frac{1}{2}}(\partial\Omega; \mathbb{C}^2) \oplus \chi_{(0, \infty)}(A)H^{-\frac{1}{2}}(\partial\Omega; \mathbb{C}^2).$$

That is,

$$u \mapsto u|_{\partial\Omega} : \text{dom}(\mathcal{D}_{\max}) \rightarrow \check{H}(A)$$

bounded surjection with kernel $\text{dom}(\mathcal{D}_{\min})$.

Boundary conditions

Boundary conditions for \mathcal{D} “live” in:

$$\check{H}(A) := \chi_{(-\infty, 0]}(A)H^{\frac{1}{2}}(\partial\Omega; \mathbb{C}^2) \oplus \chi_{(0, \infty)}(A)H^{-\frac{1}{2}}(\partial\Omega; \mathbb{C}^2).$$

That is,

$$u \mapsto u|_{\partial\Omega} : \text{dom}(\mathcal{D}_{\max}) \rightarrow \check{H}(A)$$

bounded surjection with kernel $\text{dom}(\mathcal{D}_{\min})$.

(Generalised) Atiyah-Patodi-Singer BC:

$$B_{\text{APS}} := \chi_{(-\infty, 0]}(A)H^{\frac{1}{2}}(\partial\Omega; \mathbb{C}^2).$$

Boundary conditions

Boundary conditions for \mathcal{D} “live” in:

$$\check{H}(A) := \chi_{(-\infty, 0]}(A)H^{\frac{1}{2}}(\partial\Omega; \mathbb{C}^2) \quad \oplus \quad \chi_{(0, \infty)}(A)H^{-\frac{1}{2}}(\partial\Omega; \mathbb{C}^2).$$

That is,

$$u \mapsto u|_{\partial\Omega} : \text{dom}(\mathcal{D}_{\max}) \rightarrow \check{H}(A)$$

bounded surjection with kernel $\text{dom}(\mathcal{D}_{\min})$.

(Generalised) Atiyah-Patodi-Singer BC:

$$B_{\text{APS}} := \chi_{(-\infty, 0]}(A)H^{\frac{1}{2}}(\partial\Omega; \mathbb{C}^2).$$

Eta-invariant: $\eta(A) := \eta_A(0)$ where

$$\eta_A(s) := \sum_{\lambda \in \text{spec}(A) \setminus \{0\}} \frac{\text{sgn}(\lambda)}{|\lambda|^s}.$$

Rarita-Schwinger Operator

Twisted bundle $T^*\Omega \otimes \Delta \Omega \cong \mathbb{C}^3 \otimes \mathbb{C}^2$.

Rarita-Schwinger Operator

Twisted bundle $T^*\Omega \otimes \Delta\Omega \cong \mathbb{C}^3 \otimes \mathbb{C}^2$.

Let $\iota : \Delta\Omega \rightarrow T^*\Omega \otimes \Delta\Omega$ given by

$$\iota(\psi) = -\frac{1}{3} \sum_{j=1}^3 e_j \otimes \rho(e_j)\psi.$$

Rarita-Schwinger Operator

Twisted bundle $T^*\Omega \otimes \Delta\Omega \cong \mathbb{C}^3 \otimes \mathbb{C}^2$.

Let $\iota : \Delta\Omega \rightarrow T^*\Omega \otimes \Delta\Omega$ given by

$$\iota(\psi) = -\frac{1}{3} \sum_{j=1}^3 e_j \otimes \rho(e_j)\psi.$$

$\frac{3}{2}$ -Spin bundle given by

$$\Delta_{\frac{3}{2}}\Omega := \iota(\Delta\Omega)^\perp = \ker \gamma,$$

Rarita-Schwinger Operator

Twisted bundle $T^*\Omega \otimes \Delta\Omega \cong \mathbb{C}^3 \otimes \mathbb{C}^2$.

Let $\iota : \Delta\Omega \rightarrow T^*\Omega \otimes \Delta\Omega$ given by

$$\iota(\psi) = -\frac{1}{3} \sum_{j=1}^3 e_j \otimes \rho(e_j)\psi.$$

$\frac{3}{2}$ -Spin bundle given by

$$\Delta_{\frac{3}{2}}\Omega := \iota(\Delta\Omega)^\perp = \ker \gamma,$$

where

$$\gamma(v \otimes \psi) := \rho(v)\psi.$$

Rarita-Schwinger Operator

Twisted bundle $T^*\Omega \otimes \Delta\Omega \cong \mathbb{C}^3 \otimes \mathbb{C}^2$.

Let $\iota : \Delta\Omega \rightarrow T^*\Omega \otimes \Delta\Omega$ given by

$$\iota(\psi) = -\frac{1}{3} \sum_{j=1}^3 e_j \otimes \rho(e_j)\psi.$$

$\frac{3}{2}$ -Spin bundle given by

$$\Delta_{\frac{3}{2}}\Omega := \iota(\Delta\Omega)^\perp = \ker \gamma,$$

where

$$\gamma(v \otimes \psi) := \rho(v)\psi.$$

Then,

$$T^*\Omega \otimes \Delta\Omega = \iota(\Delta\Omega) \oplus \Delta_{\frac{3}{2}}\Omega.$$

Induced Dirac operator

$$\mathcal{D}^{\text{T}^*\Omega \otimes \Delta \Omega} f = \sum_{i=1}^3 \rho(e_i)(\partial_i f) = \sum_{i=1}^3 \sum_{j,k=1}^2 (\partial_i f_{jk}) e_j \otimes \sigma_i e_k.$$

Induced Dirac operator

$$\not{D}^{\text{T}^*\Omega \otimes \not{A}\Omega} f = \sum_{i=1}^3 \rho(e_i)(\partial_i f) = \sum_{i=1}^3 \sum_{j,k=1}^2 (\partial_i f_{jk}) e_j \otimes \sigma_i e_k.$$

Orthogonal projection $\mathbf{P}_{\frac{3}{2}} : \text{T}^*\Omega \otimes \not{A}\Omega \rightarrow \not{A}_{\frac{3}{2}}\Omega$.

Induced Dirac operator

$$\not{D}^{\mathrm{T}^*\Omega \otimes \not{A}\Omega} f = \sum_{i=1}^3 \rho(e_i)(\partial_i f) = \sum_{i=1}^3 \sum_{j,k=1}^2 (\partial_i f_{jk}) e_j \otimes \sigma_i e_k.$$

Orthogonal projection $\mathbf{P}_{\frac{3}{2}} : \mathrm{T}^*\Omega \otimes \not{A}\Omega \rightarrow \not{A}_{\frac{3}{2}}\Omega$.

Rarita-Schwinger is then:

$$\not{D}_{\mathrm{RS}} := \mathbf{P}_{\frac{3}{2}} \circ \not{D}^{\mathrm{T}^*\Omega \otimes \not{A}\Omega} \Big|_{\not{A}_{\frac{3}{2}}\Omega}$$

Induced Dirac operator

$$\not{D}^{\mathrm{T}^*\Omega \otimes \not{\Delta}\Omega} f = \sum_{i=1}^3 \rho(e_i)(\partial_i f) = \sum_{i=1}^3 \sum_{j,k=1}^2 (\partial_i f_{jk}) e_j \otimes \sigma_i e_k.$$

Orthogonal projection $\mathbf{P}_{\frac{3}{2}} : \mathrm{T}^*\Omega \otimes \not{\Delta}\Omega \rightarrow \not{\Delta}_{\frac{3}{2}}\Omega$.

Rarita-Schwinger is then:

$$\not{D}_{\mathrm{RS}} := \mathbf{P}_{\frac{3}{2}} \circ \not{D}^{\mathrm{T}^*\Omega \otimes \not{\Delta}\Omega} \Big|_{\not{\Delta}_{\frac{3}{2}}\Omega}, \quad \not{D}^{\mathrm{T}^*\Omega \otimes \not{\Delta}\Omega} = \begin{pmatrix} \iota \not{D} \iota^* & \mathrm{T}^* \\ \mathrm{T} & \not{D}_{\mathrm{RS}} \end{pmatrix}.$$

Induced Dirac operator

$$\mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} f = \sum_{i=1}^3 \rho(e_i)(\partial_i f) = \sum_{i=1}^3 \sum_{j,k=1}^2 (\partial_i f_{jk}) e_j \otimes \sigma_i e_k.$$

Orthogonal projection $\mathbf{P}_{\frac{3}{2}} : T^*\Omega \otimes \mathbb{A}\Omega \rightarrow \mathbb{A}_{\frac{3}{2}}\Omega$.

Rarita-Schwinger is then:

$$\mathcal{D}_{\text{RS}} := \mathbf{P}_{\frac{3}{2}} \circ \mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} \Big|_{\mathbb{A}_{\frac{3}{2}}\Omega}, \quad \mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} = \begin{pmatrix} \iota \mathcal{D} \iota^* & T^* \\ T & \mathcal{D}_{\text{RS}} \end{pmatrix}.$$

Extract boundary adapted operator A_{RS} as before.

Induced Dirac operator

$$\mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} f = \sum_{i=1}^3 \rho(e_i)(\partial_i f) = \sum_{i=1}^3 \sum_{j,k=1}^2 (\partial_i f_{jk}) e_j \otimes \sigma_i e_k.$$

Orthogonal projection $\mathbf{P}_{\frac{3}{2}} : T^*\Omega \otimes \mathbb{A}\Omega \rightarrow \mathbb{A}_{\frac{3}{2}}\Omega$.

Rarita-Schwinger is then:

$$\mathcal{D}_{\text{RS}} := \mathbf{P}_{\frac{3}{2}} \circ \mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} \Big|_{\mathbb{A}_{\frac{3}{2}}\Omega}, \quad \mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} = \begin{pmatrix} \iota \mathcal{D} \iota^* & T^* \\ T & \mathcal{D}_{\text{RS}} \end{pmatrix}.$$

Extract boundary adapted operator A_{RS} as before.

 There is no B such that $A_{\text{RS}} + B$ is self-adjoint. 

Induced Dirac operator

$$\mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} f = \sum_{i=1}^3 \rho(e_i)(\partial_i f) = \sum_{i=1}^3 \sum_{j,k=1}^2 (\partial_i f_{jk}) e_j \otimes \sigma_i e_k.$$



Orthogonal projection $\mathbf{P}_{\frac{3}{2}} : T^*\Omega \otimes \mathbb{A}\Omega \rightarrow \mathbb{A}_{\frac{3}{2}}\Omega$.

Rarita-Schwinger is then:

$$\mathcal{D}_{\text{RS}} := \mathbf{P}_{\frac{3}{2}} \circ \mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} \Big|_{\mathbb{A}_{\frac{3}{2}}\Omega}, \quad \mathcal{D}^{T^*\Omega \otimes \mathbb{A}\Omega} = \begin{pmatrix} \iota \mathcal{D} \iota^* & T^* \\ T & \mathcal{D}_{\text{RS}} \end{pmatrix}.$$

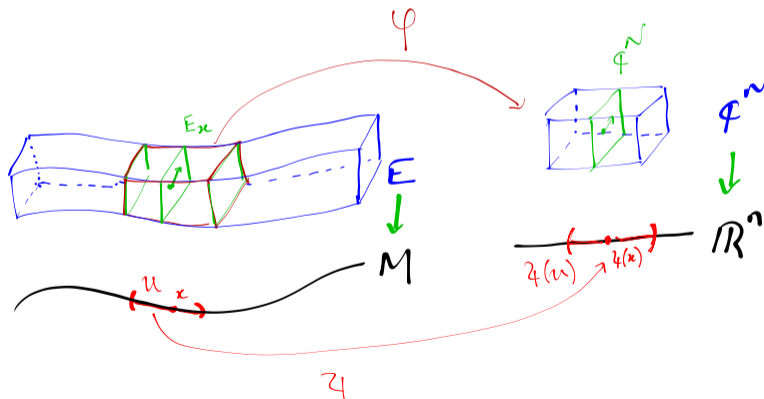
Extract boundary adapted operator A_{RS} as before.

 There is no B such that $A_{\text{RS}} + B$ is self-adjoint. 

 Self-adjointness fundamental in the Bär-Ballmann framework. 

Geometric dictionary

$(\Omega, \langle \cdot, \cdot \rangle_{\mathbb{R}^3})$	\rightsquigarrow	(M, g) manifold with metric g
$\Delta \Omega$	\rightsquigarrow	$\Delta M \rightarrow M$ spin bundle
\mathbb{C}^N	\rightsquigarrow	$E \rightarrow M$, vector bundle of rank N .



Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}})$, $(\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^{\infty}(\mathcal{E}) \rightarrow C^{\infty}(\mathcal{F})$

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^{\infty}(\mathcal{E}) \rightarrow C^{\infty}(\mathcal{F}) \rightsquigarrow \exists! D^{\dagger} : C^{\infty}(\mathcal{F}) \rightarrow C^{\infty}(\mathcal{E})$ formal adjoint.

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^{\infty}(\mathcal{E}) \rightarrow C^{\infty}(\mathcal{F}) \rightsquigarrow \exists! D^{\dagger} : C^{\infty}(\mathcal{F}) \rightarrow C^{\infty}(\mathcal{E})$ formal adjoint. I.e.,

$\forall u \in C_{cc}^{\infty}(\mathcal{E}), v \in C_{cc}^{\infty}(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^{\dagger} v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^{\infty}(\mathcal{E}) \rightarrow C^{\infty}(\mathcal{F}) \rightsquigarrow \exists! D^{\dagger} : C^{\infty}(\mathcal{F}) \rightarrow C^{\infty}(\mathcal{E})$ formal adjoint. I.e.,

$\forall u \in C_{cc}^{\infty}(\mathcal{E}), v \in C_{cc}^{\infty}(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^{\dagger} v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Define: $D_{\max} := (D^{\dagger})^*, \quad D_{\min} := \overline{D|_{C_{cc}^{\infty}(\mathcal{E})}}.$

Understand *all* closed extensions of D_{\min} sitting in D_{\max} .

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^{\infty}(\mathcal{E}) \rightarrow C^{\infty}(\mathcal{F}) \rightsquigarrow \exists! D^{\dagger} : C^{\infty}(\mathcal{F}) \rightarrow C^{\infty}(\mathcal{E})$ formal adjoint. I.e.,

$\forall u \in C_{cc}^{\infty}(\mathcal{E}), v \in C_{cc}^{\infty}(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^{\dagger} v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Define: $D_{\max} := (D^{\dagger})^*, \quad D_{\min} := \overline{D|_{C_{cc}^{\infty}(\mathcal{E})}}.$

Understand *all* closed extensions of D_{\min} sitting in D_{\max} . I.e., control

$$\text{dom}(D_{\max}) / \text{dom}(D_{\min}).$$

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^{\infty}(\mathcal{E}) \rightarrow C^{\infty}(\mathcal{F}) \rightsquigarrow \exists! D^{\dagger} : C^{\infty}(\mathcal{F}) \rightarrow C^{\infty}(\mathcal{E})$ formal adjoint. I.e.,
 $\forall u \in C_{cc}^{\infty}(\mathcal{E}), v \in C_{cc}^{\infty}(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^{\dagger}v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Define: $D_{\max} := (D^{\dagger})^*, \quad D_{\min} := \overline{D|_{C_{cc}^{\infty}(\mathcal{E})}}.$

Understand *all* closed extensions of D_{\min} sitting in D_{\max} . I.e., control

$$\text{dom}(D_{\max}) / \text{dom}(D_{\min}).$$

$\partial\mathcal{M} \neq \emptyset$ want map $\gamma : \text{dom}(D_{\max}) \rightarrow \check{H}$

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^\infty(\mathcal{E}) \rightarrow C^\infty(\mathcal{F}) \rightsquigarrow \exists! D^\dagger : C^\infty(\mathcal{F}) \rightarrow C^\infty(\mathcal{E})$ formal adjoint. I.e.,
 $\forall u \in C_{cc}^\infty(\mathcal{E}), v \in C_{cc}^\infty(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^\dagger v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Define: $D_{\max} := (D^\dagger)^*, \quad D_{\min} := \overline{D|_{C_{cc}^\infty(\mathcal{E})}}.$

Understand *all* closed extensions of D_{\min} sitting in D_{\max} . I.e., control

$$\text{dom}(D_{\max}) / \text{dom}(D_{\min}).$$

$\partial\mathcal{M} \neq \emptyset$ want map $\gamma : \text{dom}(D_{\max}) \rightarrow \check{H}$ built out of boundary trace map,

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^\infty(\mathcal{E}) \rightarrow C^\infty(\mathcal{F}) \rightsquigarrow \exists! D^\dagger : C^\infty(\mathcal{F}) \rightarrow C^\infty(\mathcal{E})$ formal adjoint. I.e.,

$\forall u \in C_{cc}^\infty(\mathcal{E}), v \in C_{cc}^\infty(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^\dagger v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Define: $D_{\max} := (D^\dagger)^*, \quad D_{\min} := \overline{D|_{C_{cc}^\infty(\mathcal{E})}}.$

Understand *all* closed extensions of D_{\min} sitting in D_{\max} . I.e., control

$$\text{dom}(D_{\max}) / \text{dom}(D_{\min}).$$

$\partial\mathcal{M} \neq \emptyset$ want map $\gamma : \text{dom}(D_{\max}) \rightarrow \check{H}$ built out of boundary trace map, bounded surjection

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^\infty(\mathcal{E}) \rightarrow C^\infty(\mathcal{F}) \rightsquigarrow \exists! D^\dagger : C^\infty(\mathcal{F}) \rightarrow C^\infty(\mathcal{E})$ formal adjoint. I.e.,
 $\forall u \in C_{cc}^\infty(\mathcal{E}), v \in C_{cc}^\infty(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^\dagger v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Define: $D_{\max} := (D^\dagger)^*, \quad D_{\min} := \overline{D|_{C_{cc}^\infty(\mathcal{E})}}.$

Understand *all* closed extensions of D_{\min} sitting in D_{\max} . I.e., control

$$\text{dom}(D_{\max}) / \text{dom}(D_{\min}).$$

$\partial\mathcal{M} \neq \emptyset$ want map $\gamma : \text{dom}(D_{\max}) \rightarrow \check{H}$ built out of boundary trace map, bounded surjection with $\ker(\gamma) = \text{dom}(D_{\min})$.

Broad aim

Hermitian vector bundles $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow (\mathcal{M}, \mu)$ meas. manifold.

Diff op $D : C^{\infty}(\mathcal{E}) \rightarrow C^{\infty}(\mathcal{F}) \rightsquigarrow \exists! D^{\dagger} : C^{\infty}(\mathcal{F}) \rightarrow C^{\infty}(\mathcal{E})$ formal adjoint. I.e.,
 $\forall u \in C_{cc}^{\infty}(\mathcal{E}), v \in C_{cc}^{\infty}(\mathcal{F}),$

$$\langle Du, v \rangle_{L^2(\mathcal{M}; \mathcal{F}, h^{\mathcal{F}})} = \langle u, D^{\dagger} v \rangle_{L^2(\mathcal{M}; \mathcal{E}, h^{\mathcal{E}})}$$

Define: $D_{\max} := (D^{\dagger})^*, \quad D_{\min} := \overline{D|_{C_{cc}^{\infty}(\mathcal{E})}}.$

Understand *all* closed extensions of D_{\min} sitting in D_{\max} . I.e., control

$$\text{dom}(D_{\max}) / \text{dom}(D_{\min}).$$

$\partial\mathcal{M} \neq \emptyset$ want map $\gamma : \text{dom}(D_{\max}) \rightarrow \check{H}$ built out of boundary trace map, bounded surjection with $\ker(\gamma) = \text{dom}(D_{\min})$. **Compute topology of \check{H} purely in terms of data on $\partial\mathcal{M}$.**

General setup

(A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;

General setup

- (A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;
- (A2) τ is an interior pointing co-vectorfield along $\partial\mathcal{M}$;

General setup

- (A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;
- (A2) τ is an interior pointing co-vectorfield along $\partial\mathcal{M}$;
- (A3) μ is a smooth volume measure on \mathcal{M} and ν is the induced smooth volume measure on $\partial\mathcal{M}$;

General setup

- (A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;
- (A2) τ is an interior pointing co-vectorfield along $\partial\mathcal{M}$;
- (A3) μ is a smooth volume measure on \mathcal{M} and ν is the induced smooth volume measure on $\partial\mathcal{M}$;
- (A4) $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow \mathcal{M}$ are Hermitian vector bundles over \mathcal{M} ;

General setup

- (A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;
- (A2) τ is an interior pointing co-vectorfield along $\partial\mathcal{M}$;
- (A3) μ is a smooth volume measure on \mathcal{M} and ν is the induced smooth volume measure on $\partial\mathcal{M}$;
- (A4) $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow \mathcal{M}$ are Hermitian vector bundles over \mathcal{M} ;
- (A5) $D : C^\infty(\mathcal{M}; \mathcal{E}) \rightarrow C^\infty(\mathcal{M}; \mathcal{F})$ is a first-order elliptic differential operator;

General setup

- (A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;
- (A2) τ is an interior pointing co-vectorfield along $\partial\mathcal{M}$;
- (A3) μ is a smooth volume measure on \mathcal{M} and ν is the induced smooth volume measure on $\partial\mathcal{M}$;
- (A4) $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow \mathcal{M}$ are Hermitian vector bundles over \mathcal{M} ;
- (A5) $D : C^\infty(\mathcal{M}; \mathcal{E}) \rightarrow C^\infty(\mathcal{M}; \mathcal{F})$ is a first-order elliptic differential operator;
- (A6) D and D^\dagger (formal adjoint of D) are complete (i.e., $C_c^\infty(\mathcal{M}; \mathcal{E})$ dense in $\text{dom}(D_{\max})$).

General setup

- (A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;
- (A2) τ is an interior pointing co-vectorfield along $\partial\mathcal{M}$;
- (A3) μ is a smooth volume measure on \mathcal{M} and ν is the induced smooth volume measure on $\partial\mathcal{M}$;
- (A4) $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow \mathcal{M}$ are Hermitian vector bundles over \mathcal{M} ;
- (A5) $D : C^\infty(\mathcal{M}; \mathcal{E}) \rightarrow C^\infty(\mathcal{M}; \mathcal{F})$ is a first-order elliptic differential operator;
- (A6) D and D^\dagger (formal adjoint of D) are complete (i.e., $C_c^\infty(\mathcal{M}; \mathcal{E})$ dense in $\text{dom}(D_{\max})$).

Consequence: reduce to cylinder $Z_{[0,T)} : [0, T) \times \partial\mathcal{M}$.

General setup

- (A1) \mathcal{M} is a manifold with compact boundary $\partial\mathcal{M} \subset \mathcal{M}$;
- (A2) τ is an interior pointing co-vectorfield along $\partial\mathcal{M}$;
- (A3) μ is a smooth volume measure on \mathcal{M} and ν is the induced smooth volume measure on $\partial\mathcal{M}$;
- (A4) $(\mathcal{E}, h^{\mathcal{E}}), (\mathcal{F}, h^{\mathcal{F}}) \rightarrow \mathcal{M}$ are Hermitian vector bundles over \mathcal{M} ;
- (A5) $D : C^\infty(\mathcal{M}; \mathcal{E}) \rightarrow C^\infty(\mathcal{M}; \mathcal{F})$ is a first-order elliptic differential operator;
- (A6) D and D^\dagger (formal adjoint of D) are complete (i.e., $C_c^\infty(\mathcal{M}; \mathcal{E})$ dense in $\text{dom}(D_{\max})$).

Consequence: reduce to cylinder $Z_{[0,T)} : [0, T) \times \partial\mathcal{M}$.

$T > 0$ determined by (A1)-(A6).

Adapted boundary operator

A adapted boundary operator to D if:

$$\sigma_A(x, \xi) = \sigma_D(x, \tau(x))^{-1} \circ \sigma_D(x, \xi).$$

Adapted boundary operator

A adapted boundary operator to D if:

$$\sigma_A(x, \xi) = \sigma_D(x, \tau(x))^{-1} \circ \sigma_D(x, \xi).$$

- Exists and are elliptic differential operators of order 1.

Adapted boundary operator

A adapted boundary operator to D if:

$$\sigma_A(x, \xi) = \sigma_D(x, \tau(x))^{-1} \circ \sigma_D(x, \xi).$$

- Exists and are elliptic differential operators of order 1.
- Unique up to an operator of order zero.

Adapted boundary operator

A adapted boundary operator to D if:

$$\sigma_A(x, \xi) = \sigma_D(x, \tau(x))^{-1} \circ \sigma_D(x, \xi).$$

- Exists and are elliptic differential operators of order 1.
- Unique up to an operator of order zero.
- Ellipticity of $D \implies$ for all $(x, \xi) \in \partial\mathcal{M} \times T^*\partial\mathcal{M}$ and $\xi \neq 0$,

$$\text{spec}(\sigma_A(x, \xi)) \cap \mathbb{R} = \emptyset.$$

Adapted boundary operator

A adapted boundary operator to D if:

$$\sigma_A(x, \xi) = \sigma_D(x, \tau(x))^{-1} \circ \sigma_D(x, \xi).$$

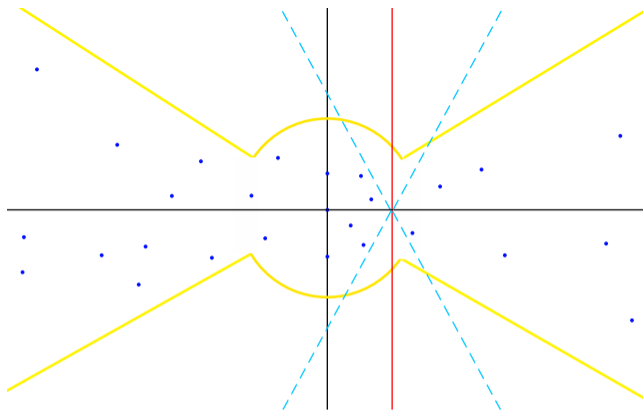
- Exists and are elliptic differential operators of order 1.
- Unique up to an operator of order zero.
- Ellipticity of $D \implies$ for all $(x, \xi) \in \partial\mathcal{M} \times T^*\partial\mathcal{M}$ and $\xi \neq 0$,

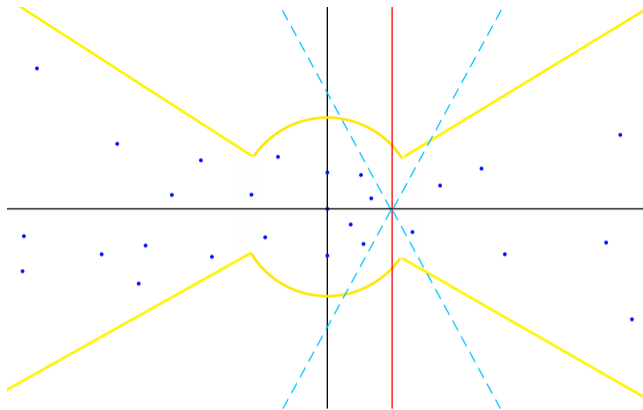
$$\text{spec}(\sigma_A(x, \xi)) \cap \mathbb{R} = \emptyset.$$

- Theorem of Shubin: there exists $\omega \in [0, \pi/2)$, $R > 0$, $C < \infty$ such that $\text{spec}(A) \subset S_\omega \cup B_R(0)$ and

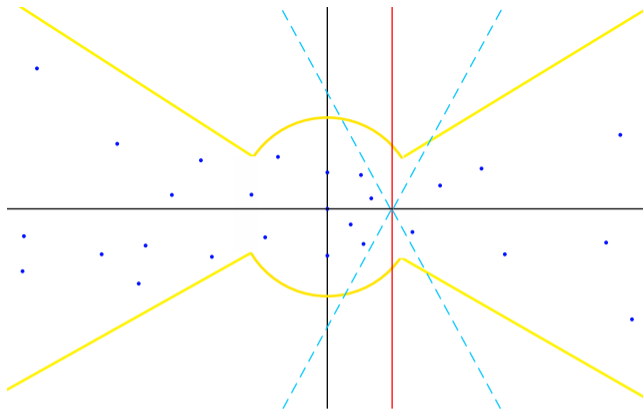
$$|\zeta| \|(\zeta - A)^{-1}\| \leq C,$$

for all $\zeta \notin S_\omega \cup B_R(0)$.

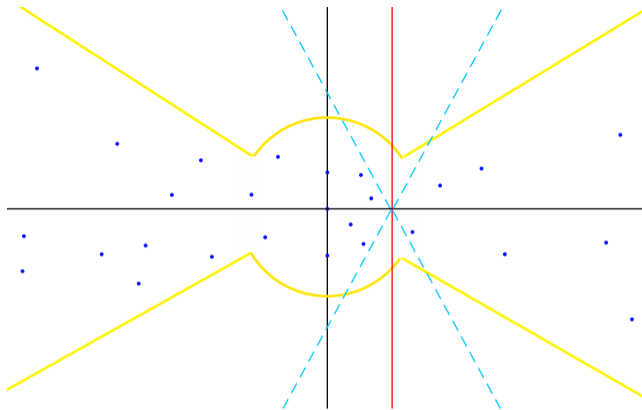




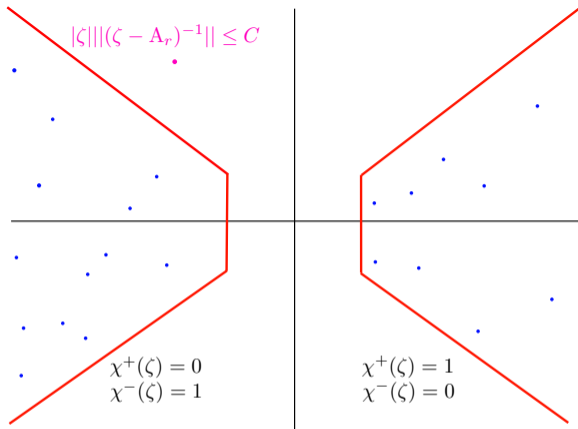
- Discrete spectrum, generally *non-orthogonal* generalised eigenspaces.

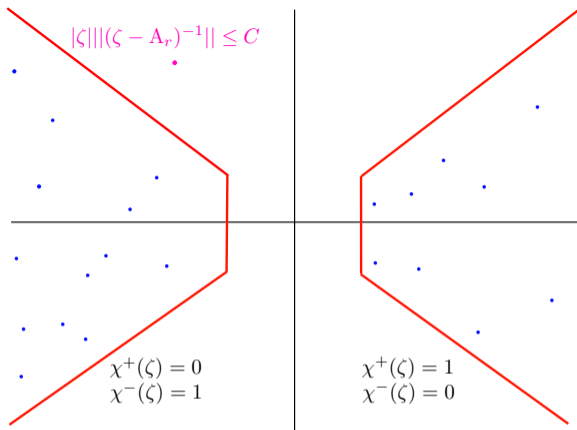


- Discrete spectrum, generally *non-orthogonal* generalised eigenspaces.
- *Admissible spectral cut* $r \in \mathbb{R}$: the line $l_r := \{\zeta \in \mathbb{C} : \operatorname{Re} \zeta = r\}$ is *not* in the spectrum of A .

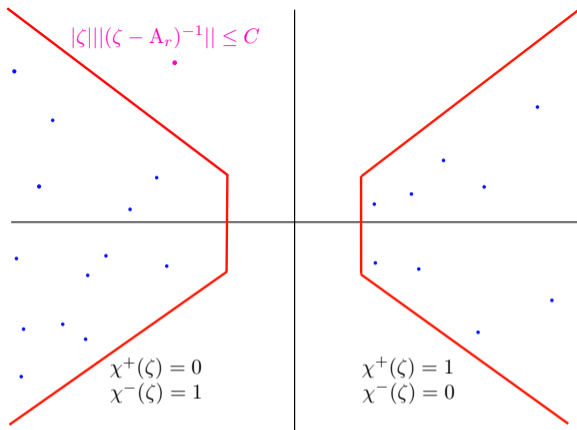


- Discrete spectrum, generally *non-orthogonal* generalised eigenspaces.
- *Admissible spectral cut* $r \in \mathbb{R}$: the line $l_r := \{\zeta \in \mathbb{C} : \operatorname{Re} \zeta = r\}$ is *not* in the spectrum of A .
- For such r , there exist $\omega_r \in [0, \pi/2)$ such that $A_r := A - r$ is invertible ω_r bi-sectorial.

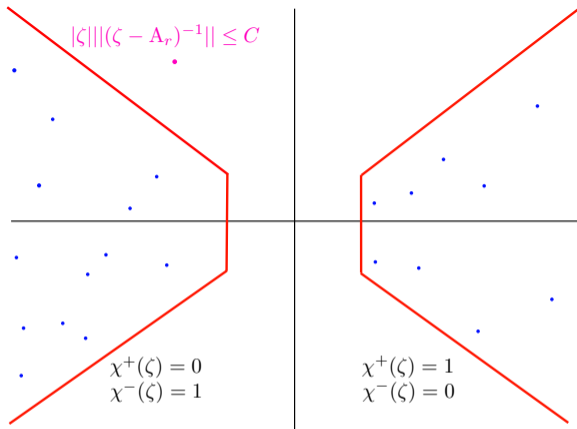




- Theorem of Grubb: $\chi^\pm(A_r)$ are Ψ DO projectors of order zero.



- Theorem of Grubb: $\chi^\pm(A_r)$ are Ψ DO projectors of order zero.
- Space: $\check{H}(A) := \chi^-(A_r)H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E}) \oplus \chi^+(A_r)H^{-\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$.



- Theorem of Grubb: $\chi^\pm(A_r)$ are Ψ DO projectors of order zero.
- Space: $\check{H}(A) := \chi^-(A_r)H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E}) \oplus \chi^+(A_r)H^{-\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$.
- Norm: $\|u\|_{\check{H}(A)}^2 := \|\chi^-(A_r)u\|_{H^{\frac{1}{2}}}^2 + \|\chi^+(A_r)u\|_{H^{-\frac{1}{2}}}^2$.

Theorem 1: Maximal domain and the $\check{H}(A)$ space

Theorem 1: Maximal domain and the $\check{H}(A)$ space

- (i) $u \mapsto u|_{\partial\mathcal{M}} : C_c^\infty(\mathcal{M}; \mathcal{E}) \rightarrow C_c^\infty(\partial\mathcal{M}; \mathcal{E})$ extends uniquely to a bounded surjection $\text{dom}(\mathbf{D}_{\max}) \rightarrow \check{H}(A)$.

Theorem 1: Maximal domain and the $\check{H}(A)$ space

(i) $u \mapsto u|_{\partial\mathcal{M}} : C_c^\infty(\mathcal{M}; \mathcal{E}) \rightarrow C_c^\infty(\partial\mathcal{M}; \mathcal{E})$ extends uniquely to a bounded surjection $\text{dom}(D_{\max}) \rightarrow \check{H}(A)$.

(ii) The space

$$\text{dom}(D_{\max}) \cap H_{\text{loc}}^1(\mathcal{M}; \mathcal{E}) = \left\{ u \in \text{dom}(D_{\max}) : u|_{\partial\mathcal{M}} \in H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E}) \right\}$$

Theorem 1: Maximal domain and the $\check{H}(A)$ space

(i) $u \mapsto u|_{\partial\mathcal{M}} : C_c^\infty(\mathcal{M}; \mathcal{E}) \rightarrow C_c^\infty(\partial\mathcal{M}; \mathcal{E})$ extends uniquely to a bounded surjection $\text{dom}(D_{\max}) \rightarrow \check{H}(A)$.

(ii) The space

$$\text{dom}(D_{\max}) \cap H_{\text{loc}}^1(\mathcal{M}; \mathcal{E}) = \left\{ u \in \text{dom}(D_{\max}) : u|_{\partial\mathcal{M}} \in H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E}) \right\}$$

(iii) The $L^2(\partial\mathcal{M}; \mathcal{E})$ inner product extends to a perfect pairing

$$\langle \cdot, \cdot \rangle : \check{H}(A) \times \check{H}(-A^*) \rightarrow \mathbb{C}.$$

(iv) For all $u \in \text{dom}(D_{\max})$ and $v \in \text{dom}(D_{\max}^\dagger)$,

$$\begin{aligned} \langle D_{\max} u, v \rangle_{L^2(\mathcal{M}; \mathcal{F})} - \langle u, D_{\max}^\dagger v \rangle_{L^2(\mathcal{M}; \mathcal{E})} \\ = - \langle u|_{\partial \mathcal{M}}, \sigma_0^* v|_{\partial \mathcal{M}} \rangle_{\check{H}(A) \times \check{H}(-A^*)}. \end{aligned}$$

(iv) For all $u \in \text{dom}(D_{\max})$ and $v \in \text{dom}(D_{\max}^\dagger)$,

$$\begin{aligned} \langle D_{\max} u, v \rangle_{L^2(\mathcal{M}; \mathcal{F})} - \langle u, D_{\max}^\dagger v \rangle_{L^2(\mathcal{M}; \mathcal{E})} \\ = - \langle u|_{\partial \mathcal{M}}, \sigma_0^* v|_{\partial \mathcal{M}} \rangle_{\check{H}(A) \times \check{H}(-A^*)}. \end{aligned}$$

(v) Higher regularity:

$$\begin{aligned} \text{dom}(D_{\max}) \cap H_{\text{loc}}^{k+1}(\mathcal{M}; \mathcal{E}) \\ = \left\{ u \in \text{dom}(D_{\max}) : Du \in H_{\text{loc}}^k(\mathcal{M}; \mathcal{F}) \right. \\ \left. \text{and } \chi^+(A_r)(u|_{\partial \mathcal{M}}) \in H^{k+\frac{1}{2}}(\partial \mathcal{M}; \mathcal{E}) \right\}. \end{aligned}$$

Boundary conditions

$B \subset \check{H}(A)$ is a *boundary condition* for D if it is a closed, linear subspace.

Boundary conditions

$B \subset \check{H}(A)$ is a *boundary condition* for D if it is a closed, linear subspace.

Associated operator domains:

$$\text{dom}(D_{B,\max}) = \{u \in \text{dom}(D_{\max}) : u|_{\partial\mathcal{M}} \in B\}$$

$$\text{dom}(D_B) = \{u \in \text{dom}(D_{\max}) \cap H_{\text{loc}}^1(\partial\mathcal{M}; \mathcal{E}) : u|_{\partial\mathcal{M}} \in B\}.$$

Boundary conditions

$B \subset \check{H}(A)$ is a *boundary condition* for D if it is a closed, linear subspace.

Associated operator domains:

$$\text{dom}(D_{B,\max}) = \{u \in \text{dom}(D_{\max}) : u|_{\partial\mathcal{M}} \in B\}$$

$$\text{dom}(D_B) = \{u \in \text{dom}(D_{\max}) \cap H_{\text{loc}}^1(\partial\mathcal{M}; \mathcal{E}) : u|_{\partial\mathcal{M}} \in B\}.$$

Similarly for the formal adjoint D^\dagger with A replaced by \tilde{A} .

- $D_{B,\max}$ closed and between D_{\min} and D_{\max} .

- $D_{B,\max}$ closed and between D_{\min} and D_{\max} .
- D_c closed extension of D_{\min} , then

$$B := \{u|_{\partial\mathcal{M}} : u \in \text{dom}(D_c)\}$$

boundary condition and $D_c = D_{B,\max}$.

- $D_{B,\max}$ closed and between D_{\min} and D_{\max} .
- D_c closed extension of D_{\min} , then

$$B := \{u|_{\partial\mathcal{M}} : u \in \text{dom}(D_c)\}$$

boundary condition and $D_c = D_{B,\max}$.

- $B \subset H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$ boundary condition if and only if $D_B = D_{B,\max}$.

- $D_{B,\max}$ closed and between D_{\min} and D_{\max} .
- D_c closed extension of D_{\min} , then

$$B := \{u|_{\partial\mathcal{M}} : u \in \text{dom}(D_c)\}$$

boundary condition and $D_c = D_{B,\max}$.

- $B \subset H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$ boundary condition if and only if $D_B = D_{B,\max}$.
- Adjoint boundary condition B^\dagger so that $(D_B)^* = D_{B^\dagger}^\dagger$:

$$B^\dagger := \left\{ v \in \check{H}(\tilde{A}) : \langle u, \sigma_0^* v \rangle_{\check{H}(A) \times \check{H}(-A^*)} = 0 \quad \forall u \in B \right\}.$$

- Classical pseudo-differential projector P of order zero (not necessarily orthogonal in L^2), the space

$$B := \overline{PH^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$$

is called a *pseudo-local boundary condition*.

- Classical pseudo-differential projector P of order zero (not necessarily orthogonal in L^2), the space

$$B := \overline{PH^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$$

is called a *pseudo-local boundary condition*.

- $B \subset H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$ a *local boundary condition* if there exists a smooth sub-bundle $E' \subset E_{\partial\mathcal{M}}$ such that

$$B = \overline{H^{\frac{1}{2}}(\partial\mathcal{M}; E')}^{\check{H}(A)}.$$

- Classical pseudo-differential projector P of order zero (not necessarily orthogonal in L^2), the space

$$B := \overline{PH^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$$

is called a *pseudo-local boundary condition*.

- $B \subset H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$ a *local boundary condition* if there exists a smooth sub-bundle $E' \subset E_{\partial\mathcal{M}}$ such that

$$B = \overline{H^{\frac{1}{2}}(\partial\mathcal{M}; E')}^{\check{H}(A)}.$$

A boundary condition B is *elliptic* if:

$$\text{dom}(D_{B,\max}) \subset H_{\text{loc}}^1(\mathcal{M}; \mathcal{E}) \quad \text{and} \quad \text{dom}(D_{B^\dagger,\max}^\dagger) \subset H_{\text{loc}}^1(\mathcal{M}; \mathcal{F})$$

Theorem 2: Pseudo-local boundary conditions

Given a pseudo-local boundary condition $B = \overline{P H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$, the following are equivalent:

Theorem 2: Pseudo-local boundary conditions

Given a pseudo-local boundary condition $B = \overline{PH^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$, the following are equivalent:

(i) B an elliptic boundary condition and $B = PH^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$,

Theorem 2: Pseudo-local boundary conditions

Given a pseudo-local boundary condition $B = \overline{PH^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$, the following are equivalent:

- (i) B an elliptic boundary condition and $B = PH^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$,
- (ii) there exists an admissible spectral cut $r \in \mathbb{R}$ and

$$P - \chi^+(A_r) : L^2(\partial\mathcal{M}; \mathcal{E}) \rightarrow L^2(\partial\mathcal{M}; \mathcal{E})$$

is Fredholm,

Theorem 2: Pseudo-local boundary conditions

Given a pseudo-local boundary condition $B = \overline{P H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$, the following are equivalent:

- (i) B an elliptic boundary condition and $B = P H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$,
- (ii) there exists an admissible spectral cut $r \in \mathbb{R}$ and

$$P - \chi^+(A_r) : L^2(\partial\mathcal{M}; \mathcal{E}) \rightarrow L^2(\partial\mathcal{M}; \mathcal{E})$$

is Fredholm,

- (iii) there exists an admissible spectral cut $r \in \mathbb{R}$ and

$$P - \chi^+(A_r) : L^2(\partial\mathcal{M}; \mathcal{E}) \rightarrow L^2(\partial\mathcal{M}; \mathcal{E})$$

is an elliptic classical pseudo of order zero.

Theorem 2: Pseudo-local boundary conditions

Given a pseudo-local boundary condition $B = \overline{P H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^{\check{H}(A)}$, the following are equivalent:

- (i) B an elliptic boundary condition and $B = P H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$,
- (ii) there exists an admissible spectral cut $r \in \mathbb{R}$ and

$$P - \chi^+(A_r) : L^2(\partial\mathcal{M}; \mathcal{E}) \rightarrow L^2(\partial\mathcal{M}; \mathcal{E})$$

is Fredholm,

- (iii) there exists an admissible spectral cut $r \in \mathbb{R}$ and

$$P - \chi^+(A_r) : L^2(\partial\mathcal{M}; \mathcal{E}) \rightarrow L^2(\partial\mathcal{M}; \mathcal{E})$$

is an elliptic classical pseudo of order zero.

In particular, if $D_B u$ is smooth, then u is smooth up to the boundary.

APS in the general setting

Given an invertible adapted boundary operator A

APS in the general setting

Given an invertible adapted boundary operator A , the boundary condition

$$B_{\text{APS}} := \chi^-(A)H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$$

is elliptic and pseudo-local.

APS in the general setting

Given an invertible adapted boundary operator A , the boundary condition

$$B_{\text{APS}} := \chi^-(A)H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$$

is elliptic and pseudo-local.

If \mathcal{M} is compact, then $D_{B_{\text{APS}}}$ is Fredholm.

APS in the general setting

Given an invertible adapted boundary operator A , the boundary condition

$$B_{\text{APS}} := \chi^-(A)H^{\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})$$

is elliptic and pseudo-local.

If \mathcal{M} is compact, then $D_{B_{\text{APS}}}$ is Fredholm.

Index formula? - Big open question.

Ingredients of the proof

Geometric reduction to the “model” operator $D_0 = \sigma_0(\partial_t + A)$:

Ingredients of the proof

Geometric reduction to the “model” operator $D_0 = \sigma_0(\partial_t + A)$:

Lemma (Lemma 4.1 (Bär-Ballmann))

On the cylinder $Z_{[0,T)}$,

$$D = \sigma_t(\partial_t + A + R_t),$$

for any adapted boundary operator A for D . The remainder term R_t is a Ψ DO of order at most one and its coefficients depend smoothly on t . Moreover,

$$\|R_t u\|_{L^2(\partial\mathcal{M};\mathcal{E})} \lesssim t \|Au\|_{L^2(\partial\mathcal{M};\mathcal{E})} + \|u\|_{L^2(\partial\mathcal{M};\mathcal{E})}$$

for $u \in C^\infty(\partial\mathcal{M};\mathcal{E})$.

Associated sectorial operators and functional calculus

- Let $\operatorname{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.

Associated sectorial operators and functional calculus

- Let $\operatorname{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.
- Define $|A_r| := A_r \operatorname{sgn}(A_r)$.

Associated sectorial operators and functional calculus

- Let $\operatorname{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.
- Define $|A_r| := A_r \operatorname{sgn}(A_r)$.
- $|A_r|$ is invertible ω_r -sectorial.

Associated sectorial operators and functional calculus

- Let $\text{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.
- Define $|A_r| := A_r \text{sgn}(A_r)$.
- $|A_r|$ is invertible ω_r -sectorial.
- Ψ DOdifferential calculus: $\text{dom}(|A_r|) = \text{dom}(|A_r|^*) = H^1(\partial\mathcal{M}; \mathcal{E})$.

Associated sectorial operators and functional calculus

- Let $\text{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.
- Define $|A_r| := A_r \text{sgn}(A_r)$.
- $|A_r|$ is invertible ω_r -sectorial.
- Ψ DO differential calculus: $\text{dom}(|A_r|) = \text{dom}(|A_r|^*) = H^1(\partial\mathcal{M}; \mathcal{E})$.
- Theorem of Auscher-McIntosh-Nahmod: $|A_r|$ has an H^∞ functional calculus.

Associated sectorial operators and functional calculus

- Let $\text{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.
- Define $|A_r| := A_r \text{sgn}(A_r)$.
- $|A_r|$ is invertible ω_r -sectorial.
- Ψ DO differential calculus: $\text{dom}(|A_r|) = \text{dom}(|A_r|^*) = H^1(\partial\mathcal{M}; \mathcal{E})$.
- Theorem of Auscher-McIntosh-Nahmod: $|A_r|$ has an H^∞ functional calculus. I.e.,

$$\int_0^\infty \|\psi(t|A_r|)u\|^2 \frac{dt}{t} \simeq \|u\|^2, \quad \forall u \in L^2(\partial\mathcal{M}; \mathcal{E})$$

Associated sectorial operators and functional calculus

- Let $\operatorname{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.
- Define $|A_r| := A_r \operatorname{sgn}(A_r)$.
- $|A_r|$ is invertible ω_r -sectorial.
- Ψ DO differential calculus: $\operatorname{dom}(|A_r|) = \operatorname{dom}(|A_r|^*) = H^1(\partial\mathcal{M}; \mathcal{E})$.
- Theorem of Auscher-McIntosh-Nahmod: $|A_r|$ has an H^∞ functional calculus. I.e.,

$$\int_0^\infty \|\psi(t|A_r|)u\|^2 \frac{dt}{t} \simeq \|u\|^2, \quad \forall u \in L^2(\partial\mathcal{M}; \mathcal{E})$$

for some (and equivalently, all) holomorphic $0 \neq \psi : S_\mu^o \rightarrow \mathbb{C}$

Associated sectorial operators and functional calculus

- Let $\operatorname{sgn}(A_r) := \chi^+(A_r) - \chi^-(A_r)$.
- Define $|A_r| := A_r \operatorname{sgn}(A_r)$.
- $|A_r|$ is invertible ω_r -sectorial.
- Ψ DO differential calculus: $\operatorname{dom}(|A_r|) = \operatorname{dom}(|A_r|^*) = H^1(\partial\mathcal{M}; \mathcal{E})$.
- Theorem of Auscher-McIntosh-Nahmod: $|A_r|$ has an H^∞ functional calculus. I.e.,

$$\int_0^\infty \|\psi(t|A_r|)u\|^2 \frac{dt}{t} \simeq \|u\|^2, \quad \forall u \in L^2(\partial\mathcal{M}; \mathcal{E})$$

for some (and equivalently, all) holomorphic $0 \neq \psi : S_\mu^o \rightarrow \mathbb{C}$ with some $C, \alpha > 0$ such that

$$|\psi(\zeta)| \leq C \min \{ |\zeta|^\alpha, |\zeta|^{-\alpha} \}.$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$.

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator
 $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator
 $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator

$\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v = e^{-t|A|}v_+ + e^{-t|A|}v_-, \quad v_\pm := \chi^\pm(A)v.$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v = e^{-t|A|}v_+ + e^{-t|A|}v_-, \quad v_\pm := \chi^\pm(A)v.$$

Show:

$$\|\mathcal{E}v\|_D^2 = \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 + \|D\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 \lesssim \|v\|_{\dot{H}(A)}^2$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator
 $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v = e^{-t|A|}v_+ + e^{-t|A|}v_-, \quad v_\pm := \chi^\pm(A)v.$$

Show:

$$\|\mathcal{E}v\|_D^2 = \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 + \|D\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 \lesssim \|v\|_{\dot{H}(A)}^2 = \|v_-\|_{H^{\frac{1}{2}}}^2 + \|v_+\|_{H^{-\frac{1}{2}}}^2.$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v = e^{-t|A|}v_+ + e^{-t|A|}v_-, \quad v_\pm := \chi^\pm(A)v.$$

Show:

$$\|\mathcal{E}v\|_D^2 = \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 + \|D\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 \lesssim \|v\|_{\dot{H}(A)}^2 = \|v_-\|_{H^{\frac{1}{2}}}^2 + \|v_+\|_{H^{-\frac{1}{2}}}^2.$$

Inhomogeneous part:

$$\|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 = \int_0^\infty \|e^{-t|A|}v\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 dt$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v = e^{-t|A|}v_+ + e^{-t|A|}v_-, \quad v_\pm := \chi^\pm(A)v.$$

Show:

$$\|\mathcal{E}v\|_D^2 = \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 + \|D\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 \lesssim \|v\|_{\dot{H}(A)}^2 = \|v_-\|_{H^{\frac{1}{2}}}^2 + \|v_+\|_{H^{-\frac{1}{2}}}^2.$$

Inhomogeneous part:

$$\begin{aligned} \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 &= \int_0^\infty \|e^{-t|A|}v\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 dt \\ &= \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{-\frac{1}{2}}v\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 \frac{dt}{t} \end{aligned}$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v = e^{-t|A|}v_+ + e^{-t|A|}v_-, \quad v_\pm := \chi^\pm(A)v.$$

Show:

$$\|\mathcal{E}v\|_D^2 = \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 + \|D\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 \lesssim \|v\|_{\dot{H}(A)}^2 = \|v_-\|_{H^{\frac{1}{2}}}^2 + \|v_+\|_{H^{-\frac{1}{2}}}^2.$$

Inhomogeneous part:

$$\begin{aligned} \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 &= \int_0^\infty \|e^{-t|A|}v\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 dt \\ &= \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{-\frac{1}{2}}v\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 \frac{dt}{t} \\ &\simeq \| |A|^{-\frac{1}{2}}v \|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 \end{aligned}$$

Role of the H^∞ calculus

Assume $M := [0, \infty) \times \partial\mathcal{M}$, $D := \sigma_0(\partial_t + A)$. Extension operator $\mathcal{E} : C_c^\infty(\partial\mathcal{M}; \mathcal{E}) \rightarrow \text{dom}(D_{\max})$

$$\mathcal{E}v := e^{-t|A|}v = e^{-t|A|}v_+ + e^{-t|A|}v_-, \quad v_\pm := \chi^\pm(A)v.$$

Show:

$$\|\mathcal{E}v\|_D^2 = \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 + \|D\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 \lesssim \|v\|_{\dot{H}(A)}^2 = \|v_-\|_{H^{\frac{1}{2}}}^2 + \|v_+\|_{H^{-\frac{1}{2}}}^2.$$

Inhomogeneous part:

$$\begin{aligned} \|\mathcal{E}v\|_{L^2(\mathcal{M}; \mathcal{E})}^2 &= \int_0^\infty \|e^{-t|A|}v\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 dt \\ &= \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{-\frac{1}{2}}v\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 \frac{dt}{t} \\ &\simeq \| |A|^{-\frac{1}{2}}v \|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 \simeq \|v\|_{H^{-\frac{1}{2}}(\partial\mathcal{M}; \mathcal{E})}^2. \end{aligned}$$

Firstly,

$$D\mathcal{E}v_- = \sigma_0(\partial_t + A)\mathcal{E}v_-$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \end{aligned}$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \end{aligned}$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \\ &= -2\sigma_0|A|e^{-t|A|}v_-. \end{aligned}$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \\ &= -2\sigma_0|A|e^{-t|A|}v_-. \end{aligned}$$

Then,

$$\|D\mathcal{E}v_-\|_{L^2(\mathcal{M};\mathcal{E})}^2 = 4 \int_0^\infty \|\sigma_0|A|e^{-t|A|}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 dt$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \\ &= -2\sigma_0|A|e^{-t|A|}v_-. \end{aligned}$$

Then,

$$\begin{aligned} \|D\mathcal{E}v_-\|_{L^2(\mathcal{M};\mathcal{E})}^2 &= 4 \int_0^\infty \|\sigma_0|A|e^{-t|A|}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 dt \\ &\simeq \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \frac{dt}{t} \end{aligned}$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \\ &= -2\sigma_0|A|e^{-t|A|}v_-. \end{aligned}$$

Then,

$$\begin{aligned} \|D\mathcal{E}v_-\|_{L^2(\mathcal{M};\mathcal{E})}^2 &= 4 \int_0^\infty \|\sigma_0|A|e^{-t|A|}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 dt \\ &\simeq \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \frac{dt}{t} \\ &\simeq \| |A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \end{aligned}$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \\ &= -2\sigma_0|A|e^{-t|A|}v_-. \end{aligned}$$

Then,

$$\begin{aligned} \|D\mathcal{E}v_-\|_{L^2(\mathcal{M};\mathcal{E})}^2 &= 4 \int_0^\infty \|\sigma_0|A|e^{-t|A|}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 dt \\ &\simeq \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \frac{dt}{t} \\ &\simeq \| |A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \simeq \|v_-\|_{H^{\frac{1}{2}}(\partial\mathcal{M};\mathcal{E})}^2. \end{aligned}$$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \\ &= -2\sigma_0|A|e^{-t|A|}v_-. \end{aligned}$$

Then,

$$\begin{aligned} \|D\mathcal{E}v_-\|_{L^2(\mathcal{M};\mathcal{E})}^2 &= 4 \int_0^\infty \|\sigma_0|A|e^{-t|A|}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 dt \\ &\simeq \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \frac{dt}{t} \\ &\simeq \| |A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \simeq \|v_-\|_{H^{\frac{1}{2}}(\partial\mathcal{M};\mathcal{E})}^2. \end{aligned}$$

Combining with $D\mathcal{E}v_+ = \sigma_0(\partial_t + A)v_+ = 0$

Firstly,

$$\begin{aligned} D\mathcal{E}v_- &= \sigma_0(\partial_t + A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t + |A|\operatorname{sgn}(A))\chi^-(A)\mathcal{E}v_- \\ &= \sigma_0(\partial_t - |A|)e^{-t|A|}v_- \\ &= -2\sigma_0|A|e^{-t|A|}v_-. \end{aligned}$$

Then,

$$\begin{aligned} \|D\mathcal{E}v_-\|_{L^2(\mathcal{M};\mathcal{E})}^2 &= 4 \int_0^\infty \|\sigma_0|A|e^{-t|A|}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 dt \\ &\simeq \int_0^\infty \|t^{\frac{1}{2}}|A|^{\frac{1}{2}}e^{-t|A|}|A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \frac{dt}{t} \\ &\simeq \| |A|^{\frac{1}{2}}v_-\|_{L^2(\partial\mathcal{M};\mathcal{E})}^2 \simeq \|v_-\|_{\mathbb{H}^{\frac{1}{2}}(\partial\mathcal{M};\mathcal{E})}^2. \end{aligned}$$

Combining with $D\mathcal{E}v_+ = \sigma_0(\partial_t + A)v_+ = 0$, obtain:

$$\|\mathcal{E}v\|_D^2 \lesssim \|v_-\|_{\mathbb{H}^{\frac{1}{2}}(\partial\mathcal{M};\mathcal{E})}^2 + \|v_+\|_{\mathbb{H}^{-\frac{1}{2}}(\partial\mathcal{M};\mathcal{E})}^2 = \|v\|_{\check{\mathbb{H}}(A)}^2.$$

Maximal regularity

Banach-valued Cauchy problem: $f \in L^2(Z_{[0,\rho]}; \mathcal{E})$,

$$\partial_t W(t; f) + |A_r|W(t; f) = f(t)$$

$$\lim_{t \rightarrow 0} W(t; f) = 0.$$

Maximal regularity

Banach-valued Cauchy problem: $f \in L^2(Z_{[0,\rho]}; \mathcal{E})$,

$$\partial_t W(t; f) + |A_r| W(t; f) = f(t)$$

$$\lim_{t \rightarrow 0} W(t; f) = 0.$$

Solution given by:

$$W(t; f) = \int_0^t e^{-(t-s)|A_r|} f(s) ds.$$

Maximal regularity

Banach-valued Cauchy problem: $f \in L^2(Z_{[0,\rho]}; \mathcal{E})$,

$$\partial_t W(t; f) + |A_r|W(t; f) = f(t)$$

$$\lim_{t \rightarrow 0} W(t; f) = 0.$$

Solution given by:

$$W(t; f) = \int_0^t e^{-(t-s)|A_r|} f(s) ds.$$

Key estimate - maximal regularity:

$$\int_0^\rho \|\partial_t W(t; f)\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 dt + \int_0^\rho \| |A_r| W(t; f) \|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 dt \lesssim \int_0^\rho \|f(t)\|_{L^2(\partial\mathcal{M}; \mathcal{E})}^2 dt.$$

Define:

$$S_{0,r}u(t) = \int_0^t e^{-(t-s)|A_r|} \chi^+(A_r) u(s) ds \\ - \int_t^\rho e^{-(s-t)|A_r|} \chi^-(A_r) u(s) ds$$

Define:

$$S_{0,r}u(t) = \int_0^t e^{-(t-s)|A_r|} \chi^+(A_r) u(s) ds \\ - \int_t^\rho e^{-(s-t)|A_r|} \chi^-(A_r) u(s) ds$$

Let $(C_\rho u)(s) = u(\rho - s)$,

(i) $(\partial_t + A_r) S_{0,r} = I$.

Define:

$$S_{0,r}u(t) = \int_0^t e^{-(t-s)|A_r|} \chi^+(A_r) u(s) ds \\ - \int_t^\rho e^{-(s-t)|A_r|} \chi^-(A_r) u(s) ds$$

Let $(C_\rho u)(s) = u(\rho - s)$,

- (i) $(\partial_t + A_r) S_{0,r} = I$.
- (ii) $S_{0,r} : H^k(Z_{[0,\rho]}; \mathcal{E}) \rightarrow H^{k+1}(Z_{[0,\rho]}; \mathcal{E})$ bounded.

Define:

$$S_{0,r}u(t) = \int_0^t e^{-(t-s)|A_r|} \chi^+(A_r) u(s) ds \\ - \int_t^\rho e^{-(s-t)|A_r|} \chi^-(A_r) u(s) ds$$

Let $(C_\rho u)(s) = u(\rho - s)$,

- (i) $(\partial_t + A_r) S_{0,r} = I$.
- (ii) $S_{0,r} : H^k(Z_{[0,\rho]}; \mathcal{E}) \rightarrow H^{k+1}(Z_{[0,\rho]}; \mathcal{E})$ bounded.
- (iii) whenever $u(\rho) = 0$ (or $\text{spt } u \subset Z_{[0,\rho)}$),

Define:

$$S_{0,r}u(t) = \int_0^t e^{-(t-s)|A_r|} \chi^+(A_r) u(s) ds \\ - \int_t^\rho e^{-(s-t)|A_r|} \chi^-(A_r) u(s) ds$$

Let $(C_\rho u)(s) = u(\rho - s)$,

- (i) $(\partial_t + A_r) S_{0,r} = I$.
- (ii) $S_{0,r} : H^k(Z_{[0,\rho]}; \mathcal{E}) \rightarrow H^{k+1}(Z_{[0,\rho]}; \mathcal{E})$ bounded.
- (iii) whenever $u(\rho) = 0$ (or $\text{spt } u \subset Z_{[0,\rho)}$),

$$(I - S_{0,r} (\partial_t + A_r)) u = e^{-t|A_r|} (\chi^+(A_r) u(0)).$$

Future program

Using current viewpoint as a *template*

Future program

Using current viewpoint as a *template*:

- *General order case* [Magnus Goffeng (Lund), Hemanth Saratchandran (Adelaide)]

Future program

Using current viewpoint as a *template*:

- *General order case* [Magnus Goffeng (Lund), Hemanth Saratchandran (Adelaide)]
Seeley (1965) gives a “Czech” space: mixed-order Sobolev spaces via Calderón projectors.

Future program

Using current viewpoint as a *template*:

- *General order case* [Magnus Goffeng (Lund), Hemanth Saratchandran (Adelaide)]
Seeley (1965) gives a “Czech” space: mixed-order Sobolev spaces via Calderón projectors.
- *Lipschitz boundary* [Andreas Rosén (Gothenburg) and Magnus Goffeng (Lund)]

Future program

Using current viewpoint as a *template*:

- *General order case* [Magnus Goffeng (Lund), Hemanth Saratchandran (Adelaide)]
Seeley (1965) gives a “Czech” space: mixed-order Sobolev spaces via Calderón projectors.
- *Lipschitz boundary* [Andreas Rosén (Gothenburg) and Magnus Goffeng (Lund)]
Quadratic estimates to be proved directly - methods from the Kato square root problem:

Future program

Using current viewpoint as a *template*:

- *General order case* [Magnus Goffeng (Lund), Hemanth Saratchandran (Adelaide)]
Seeley (1965) gives a “Czech” space: mixed-order Sobolev spaces via Calderón projectors.
- *Lipschitz boundary* [Andreas Rosén (Gothenburg) and Magnus Goffeng (Lund)]
Quadratic estimates to be proved directly - methods from the Kato square root problem:
dyadic decomposition,

Future program

Using current viewpoint as a *template*:

- *General order case* [Magnus Goffeng (Lund), Hemanth Saratchandran (Adelaide)]
Seeley (1965) gives a “Czech” space: mixed-order Sobolev spaces via Calderón projectors.
- *Lipschitz boundary* [Andreas Rosén (Gothenburg) and Magnus Goffeng (Lund)]
Quadratic estimates to be proved directly - methods from the Kato square root problem:
dyadic decomposition, off-diagonal estimates (automatic for first-order),

Future program

Using current viewpoint as a *template*:

- *General order case* [Magnus Goffeng (Lund), Hemanth Saratchandran (Adelaide)]
Seeley (1965) gives a “Czech” space: mixed-order Sobolev spaces via Calderón projectors.
- *Lipschitz boundary* [Andreas Rosén (Gothenburg) and Magnus Goffeng (Lund)]
Quadratic estimates to be proved directly - methods from the Kato square root problem:
dyadic decomposition, off-diagonal estimates (automatic for first-order), reduce to local $T(b)$ theorem and Carleson measure estimate.

- *Nonlinear problems - L^p estimates*

- *Nonlinear problems - L^p estimates*

$$\|\mathcal{E}u\|_{L^p(\mathcal{M};\mathcal{E})}^p = \int_0^\infty \int_{\partial\mathcal{M}} |e^{-t|A|}u|^p d\mu_{\partial\mathcal{M}} dt = \int_0^\infty \|t^{\frac{1}{p}}e^{-t|A|}u\|^p \frac{dt}{t}.$$

- *Nonlinear problems - L^p estimates*

$$\|\mathcal{E}u\|_{L^p(\mathcal{M};\mathcal{E})}^p = \int_0^\infty \int_{\partial\mathcal{M}} |e^{-t|A|}u|^p d\mu_{\partial\mathcal{M}} dt = \int_0^\infty \|t^{\frac{1}{p}}e^{-t|A|}u\|^p \frac{dt}{t}.$$

Leads to: Besov space data on the boundary.

- *Nonlinear problems - L^p estimates*

$$\|\mathcal{E}u\|_{L^p(\mathcal{M};\mathcal{E})}^p = \int_0^\infty \int_{\partial\mathcal{M}} |e^{-t|A|}u|^p d\mu_{\partial\mathcal{M}} dt = \int_0^\infty \|t^{\frac{1}{p}}e^{-t|A|}u\|^p \frac{dt}{t}.$$

Leads to: Besov space data on the boundary.

Guess: $\check{H}_p := \chi^-(A)\mathbb{B}_{p,p}^{1-\frac{1}{p}}(\partial\mathcal{M};\mathcal{E}) \oplus \chi^+(A)\mathbb{B}_{p,p}^{-\frac{1}{p}}(\partial\mathcal{M};E).$

- *η -invariants for non-Dirac type operators*

- *Nonlinear problems - L^p estimates*

$$\|\mathcal{E}u\|_{L^p(\mathcal{M};\mathcal{E})}^p = \int_0^\infty \int_{\partial\mathcal{M}} |e^{-t|A|}u|^p d\mu_{\partial\mathcal{M}} dt = \int_0^\infty \|t^{\frac{1}{p}}e^{-t|A|}u\|^p \frac{dt}{t}.$$

Leads to: Besov space data on the boundary.

Guess: $\check{H}_p := \chi^-(A)\mathbb{B}_{p,p}^{1-\frac{1}{p}}(\partial\mathcal{M};\mathcal{E}) \oplus \chi^+(A)\mathbb{B}_{p,p}^{-\frac{1}{p}}(\partial\mathcal{M};E).$

- *η -invariants for non-Dirac type operators*

Recall:

$$\eta_A(s) := \sum_{\lambda \in \text{spec}(A) \setminus \{0\}} \frac{\text{sgn}(\lambda)}{|\lambda|^s}.$$

- *Nonlinear problems - L^p estimates*

$$\|\mathcal{E}u\|_{L^p(\mathcal{M};\mathcal{E})}^p = \int_0^\infty \int_{\partial\mathcal{M}} |e^{-t|A|}u|^p d\mu_{\partial\mathcal{M}} dt = \int_0^\infty \|t^{\frac{1}{p}}e^{-t|A|}u\|^p \frac{dt}{t}.$$

Leads to: Besov space data on the boundary.

Guess: $\check{H}_p := \chi^-(A)\mathbb{B}_{p,p}^{1-\frac{1}{p}}(\partial\mathcal{M};\mathcal{E}) \oplus \chi^+(A)\mathbb{B}_{p,p}^{-\frac{1}{p}}(\partial\mathcal{M};E).$

- *η -invariants for non-Dirac type operators*

Recall:

$$\eta_A(s) := \sum_{\lambda \in \text{spec}(A) \setminus \{0\}} \frac{\text{sgn}(\lambda)}{|\lambda|^s}.$$

$\eta_A(0)$ defined?

- *Nonlinear problems - L^p estimates*

$$\|\mathcal{E}u\|_{L^p(\mathcal{M};\mathcal{E})}^p = \int_0^\infty \int_{\partial\mathcal{M}} |e^{-t|A|}u|^p d\mu_{\partial\mathcal{M}} dt = \int_0^\infty \|t^{\frac{1}{p}}e^{-t|A|}u\|^p \frac{dt}{t}.$$

Leads to: Besov space data on the boundary.

Guess: $\check{H}_p := \chi^-(A)\mathbb{B}_{p,p}^{1-\frac{1}{p}}(\partial\mathcal{M};\mathcal{E}) \oplus \chi^+(A)\mathbb{B}_{p,p}^{-\frac{1}{p}}(\partial\mathcal{M};E).$

- *η -invariants for non-Dirac type operators*

Recall:

$$\eta_A(s) := \sum_{\lambda \in \text{spec}(A) \setminus \{0\}} \frac{\text{sgn}(\lambda)}{|\lambda|^s}.$$

$\eta_A(0)$ defined? $s \mapsto \eta_A(s)$ analytic?

- *Nonlinear problems - L^p estimates*

$$\|\mathcal{E}u\|_{L^p(\mathcal{M};\mathcal{E})}^p = \int_0^\infty \int_{\partial\mathcal{M}} |e^{-t|A|}u|^p d\mu_{\partial\mathcal{M}} dt = \int_0^\infty \|t^{\frac{1}{p}}e^{-t|A|}u\|^p \frac{dt}{t}.$$

Leads to: Besov space data on the boundary.

Guess: $\check{H}_p := \chi^-(A)\mathbb{B}_{p,p}^{1-\frac{1}{p}}(\partial\mathcal{M};\mathcal{E}) \oplus \chi^+(A)\mathbb{B}_{p,p}^{-\frac{1}{p}}(\partial\mathcal{M};E).$

- *η -invariants for non-Dirac type operators*

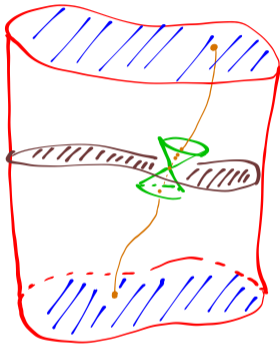
Recall:

$$\eta_A(s) := \sum_{\lambda \in \text{spec}(A) \setminus \{0\}} \frac{\text{sgn}(\lambda)}{|\lambda|^s}.$$

$\eta_A(0)$ defined? $s \mapsto \eta_A(s)$ analytic?

H^∞ -functional calculus and harmonic analysis: alternative perspective of Atiyah-Patodi-Singer.

- Lorentzian manifolds with spacelike boundary

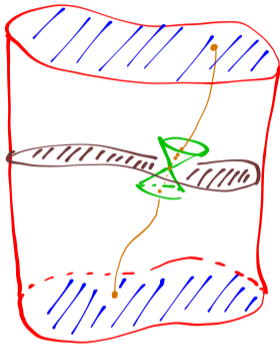


$$(\Sigma_1, g_1)$$

$$M = [0, 1] \times \Sigma$$
$$g = -N^2 dt + g_t$$

$$(\Sigma_0, g_0)$$

- Lorentzian manifolds with spacelike boundary



$$(\Sigma_1, g_1)$$

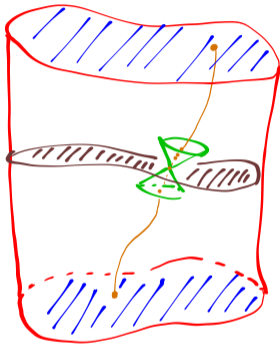
$$M = [0, 1] \times \Sigma$$

$$g = -N^2 dt + g_t$$

$$(\Sigma_0, g_0)$$

Extension operator: wave propagation operator.

- Lorentzian manifolds with spacelike boundary



$$(\Sigma_1, g_1)$$

$$M = [0, 1] \times \Sigma$$

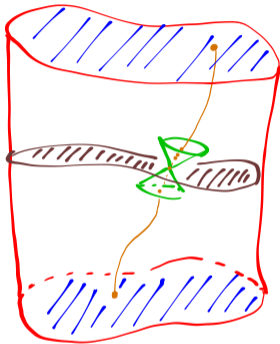
$$g = -N^2 dt + g_t$$

$$(\Sigma_0, g_0)$$

Extension operator: wave propagation operator.

Bisectoriality is a problem: need strip type or similar.

- Lorentzian manifolds with spacelike boundary



$$(\Sigma_1, g_1)$$

$$M = [0, 1] \times \Sigma$$

$$g = -N^2 dt + g_t$$

$$(\Sigma_0, g_0)$$

Extension operator: wave propagation operator.

Bisectoriality is a problem: need strip type or similar.

Key idea: identify the right function spaces.