# $L^{p}$-resolvent estimates for compact Riemannian manifolds and hyperbolic space 

Chris Sogge (Johns Hopkins University)
Joint work with:

Jean Bourgain (IAS)<br>Peng Shao (SUNY, Binghamton)<br>Xiaohua Yao (Central China Normal U)

and Shanlin Huang (Huazhong Univ and JHU)

## I.) Basic problems

Interested in uniform estimates of the form

$$
\|u\|_{L^{s}(M)} \leq C\|(\Delta+\zeta) u\|_{L^{r}(M)}, \quad u \in C_{0}^{\infty}, \quad \zeta \in \mathcal{R} \subset \mathbb{C}
$$

- $\Delta=$ Laplacian on ( $M, g$ ), norms taken w.r.t. volume element
- $(M, g)$ either compact w/out boundary, or simply connected with constant curvature (e.g., $\mathbb{H}^{n}$ )
- $M$ compact $\zeta=\lambda+i$ for appropriate $C=C(\lambda)=\lambda^{\sigma(r, s)}$ implies bounds for eigenfunctions
- $M$ noncompact and $r=s=2$, cutoff resolvent bounds for $(\Delta+\zeta)^{-1}$ important in semiclassical analysis (resonances) and sensitive to trapping of geodesics-distribution of resonances.
- Applications to control theory.
- For $L^{p}$-resolvent bounds and $M$ compact or noncompact: Expect estimates to be related to spectrum of $\Delta$ and dynamics of geodesic flow


## II a.) Optimal results for Euclidean space

Suppose that

- $n\left(\frac{1}{r}-\frac{1}{s}\right)=2$
- and $\min \left(\left|\frac{1}{r}-\frac{1}{2}\right|,\left|\frac{1}{s}-\frac{1}{2}\right|\right)>\frac{1}{2 n}$, (i.e., $\frac{2 n}{n-1}<s<\frac{2 n}{n-3}$ )

Kenig, Ruiz, S. ('87):
If $n \geq 3$ and $r, s$ as above $\exists C_{r, s}$ so that

$$
\begin{equation*}
\|u\|_{L^{s}\left(\mathbb{R}^{n}\right)} \leq C_{r, s}\|(\Delta+\zeta) u\|_{L^{r}\left(\mathbb{R}^{n}\right)}, \quad \zeta \in \mathbb{C}, u \in C_{0}^{\infty} . \tag{1}
\end{equation*}
$$

Idea: If $f=(\Delta+\zeta) u$, then

$$
\begin{gathered}
u(x)=(2 \pi)^{-n} \int_{\mathbb{R}^{n}} e^{i x \cdot \xi} \frac{\hat{f}(\xi)}{\zeta-|\xi|^{2}} d \xi \\
\hat{f}(\xi)=\int_{\mathbb{R}^{n}} e^{-i y \cdot \xi} f(y) d y, \quad f(x)=(2 \pi)^{-n} \int_{\mathbb{R}^{n}} e^{i x \cdot \xi} \hat{f}(\xi) d \xi
\end{gathered}
$$

## II b.) Euclidean restriction estimates

Worse case: $\zeta=1-i \varepsilon, \varepsilon \searrow 0$, since

$$
\operatorname{Im}\left(1-i \varepsilon-|\xi|^{2}\right)^{-1} \rightarrow \pi d S,\left(\text { surface measure on } S^{n-1}\right)
$$

So (1) implies that

$$
\begin{equation*}
\left\|\int_{S^{n-1}} e^{i x \cdot \omega} \hat{f}(\omega) d S(\omega)\right\|_{L^{s}\left(\mathbb{R}^{n}\right)} \leq C\|f\|_{L^{r}\left(\mathbb{R}^{n}\right)}, \tag{2}
\end{equation*}
$$

- See that you need $s>\frac{2 n}{n-1}$ if $f \in \mathcal{S}$ with $\hat{f}=1$ near $S^{n-1}$
- Case where $s^{\prime}=r=\frac{2 n}{n-2} \Longleftrightarrow$ to special case of Stein-Tomas:

$$
\begin{equation*}
\left(\int_{S^{n-1}}|\hat{u}|^{2} d S\right)^{\frac{1}{2}} \leq C\|u\|_{L^{p}\left(\mathbb{R}^{n}\right)}, \quad 1 \leq p \leq \frac{2(n+1)}{n+3} \tag{3}
\end{equation*}
$$

KRS: Can reverse argument that $(1) \Longrightarrow(2)$ as well.

## Remarks

If $d E_{\lambda}, \lambda \in[0, \infty)$ denotes the spectral measure for $\sqrt{-\Delta_{\mathbb{R}^{n}}}$, then, after rescaling, (2) is equivalent to

$$
\begin{align*}
\left\|d E_{\lambda}\right\|_{L^{r}\left(\mathbb{R}^{n}\right) \rightarrow L^{s}\left(\mathbb{R}^{n}\right)}=\lim _{\varepsilon \searrow 0} \varepsilon^{-1}\left\|1_{[\lambda, \lambda+\varepsilon]}\left(\sqrt{-\Delta_{\mathbb{R}^{n}}}\right)\right\|_{L^{r} \rightarrow L^{s}} & \\
& \leq C \lambda \tag{4}
\end{align*}
$$

- Continuous spectrum of $\Delta_{\mathbb{R}^{n}}$ responsible for this as well as resolvent Euclidean estimates (1).
- Expect different story for compact manifolds, and, there, expect dynamics of geodesic flow and spectral properties of $\Delta_{g}$ to dictate how close you can come to (1) and (4)
- Expect also: Resolvent bounds corresponding to high frequencies can see spectrum and dynamics. Known to be true in many cases for $L^{2}$ —also for $L^{p}$ bounds.
- What about $\mathbb{H}^{n}$ ?


## III.) Compact manifolds: Classical Spectral Theory

In 1910 Sommerfeld, followed 3 months later by Lorentz, gave famous lectures inspiring Weyl's work. Sommerfeld interested in forced vibration problem in dimensions $n=1,2,3:(\Delta+\zeta) u(x)=f(x)$,
$x \in \Omega \Subset \mathbb{R}^{n},\left.u\right|_{\partial \Omega}=0$
Asked how properties of solution operator $(\Delta+\zeta)^{-1}$ related to solutions of the free vibration problem

$$
\left(\Delta+\lambda_{j}^{2}\right) e_{j}(x)=0,\left.\quad e_{j}\left|\partial \Omega=0, \quad \int_{\Omega}\right| e_{j}\right|^{2} d x=1
$$

Kernel of solution operator: $\mathfrak{S}(x, y)=\sum_{j} \frac{e_{j}(x) e_{j}(y)}{\zeta-\lambda_{j}^{2}}$ Sommerfeld reasoned that $(\Delta+\zeta) \mathfrak{S}(x, y)=\sum e_{j}(x) e_{j}(y)$ is "spike function"

Also conjectured that properties of $\mathfrak{S}(x, y)$ should be related to distribution of eigenvalues $\left\{\lambda_{j}\right\}$, and "cancellation from numerator" (oscillation of e.f.'s)

## Weyl Law

Lorentz's subsequent 1910 lecture spelled out the eigenvalue problem more precisely and asked whether for the eigenvalues for the Dirichlet Laplacian in smooth domains $\Omega \subset \mathbb{R}^{n}$ one has for $N(\lambda)=$ number $\lambda_{j} \leq \lambda$

$$
N(\lambda)=(2 \pi)^{-n}(\operatorname{Vol} B)(\operatorname{Vol} \Omega) \lambda^{n}+o\left(\lambda^{n}\right)
$$

Hilbert: No way in my lifetime

Weyl: Yes! (several proofs in 1911-12 comparison arguments, heat kernel, Tauberian arguments...)

Sharp Weyl formula (Avakumovic '50s):

$$
N(\lambda)=c_{M} \lambda^{n}+O\left(\lambda^{n-1}\right)
$$

Improvements over the years in many special cases with "good" dynamics

## IV.) Resolvent bounds for compact manifolds

Given a compact Riemannian manifold ( $M, g$ ) of dimension $n \geq 3$, interested in regions $\mathcal{R}(g)$ for which one can have uniform resolvent estimates:

$$
\|u\|_{L^{\frac{2 n}{n-2}}(M)} \leq C\left\|\left(\Delta_{g}+\zeta\right) u\right\|_{L^{\frac{2 n}{n+2}(M)}}, \zeta \in \mathcal{R}(g), u \in C^{\infty} .
$$

Z. Shen (2001): For the torus $\mathbb{T}^{\mathbf{n}}=\mathbb{R}^{\mathbf{n}} / \mathbb{Z}^{\mathbf{n}}$ can take region to be

$$
\left\{\zeta \in \mathbb{C}: \operatorname{Re} \zeta \leq(\operatorname{Im} \zeta)^{2},|\zeta| \geq 1\right\}
$$

Dos Santos Ferreira, Kenig and Salo [DKS] (2011): Same results for any compact Riemannian manifold

## Problem raised by DKS for $\left(\Delta_{g}+\zeta\right)^{-1}: L^{\frac{2 n}{n+2}} \rightarrow L^{\frac{2 n}{n-2}}$

Does the DKS-S theorem hold for a larger region, specifically the region outside the curve $\gamma_{\text {opt }}$, which is $|\operatorname{lm} \zeta|=1$ (unit distance from spectrum of $-\Delta_{g}$ )?
This would be natural Riemannian version of KRS results for $\mathbb{R}^{n}$.


## V.) Weyl law \& answer (w Bourgain, Shao \& Yao [BSSY])

The answer is NO: In some cases cannot come close to $\gamma_{o p t}$ and earlier bounds of DKS in fact cannot be improved:
Write $\zeta=(\lambda+i \varepsilon(\lambda))^{2}=\lambda^{2}+2 i \varepsilon \lambda-\varepsilon^{2}$.
Then $\varepsilon(\lambda)=1$ corresponds to $\gamma_{D K S S}$ curve. (BSSY): If
$\left(\Delta+(\lambda+i \varepsilon(\lambda))^{2}\right)^{-1}: L^{\frac{2 n}{n+2}} \rightarrow L^{\frac{2 n}{n-2}}$ uniformly, then

$$
\begin{aligned}
& \#\left\{\lambda_{j}: \lambda_{j} \in[\lambda-\varepsilon(\lambda), \lambda+\varepsilon(\lambda)]\right\} \\
& \leq \operatorname{Vol}\left\{\xi \in \mathbb{R}^{n}:|\xi| \in[\lambda-\varepsilon(\lambda), \lambda+\varepsilon(\lambda)]\right\} \lesssim \varepsilon(\lambda) \lambda^{n-1} \\
& \quad=\text { Volume of } \varepsilon(\lambda)-\text { annulus about } \lambda S^{n-1} .
\end{aligned}
$$

Cannot hold if $\varepsilon(\lambda) \searrow 0$ for $S^{n}$ as distinct eigenvalues are $\lambda=\sqrt{k(k+n-1)}$, repeating with multiplicity $\approx \lambda^{n-1}$.
$\gamma_{\text {opt }}$ corresponds to $e(\lambda)=\lambda^{-1}$. DIFFICULT!!
Sommerfeld correct: Resolvent operators sensitive to spectral properties of $\Delta_{g}$

## Some cases of nonclustering spectrum

Some manifolds $(M, g)$ where it is known that $\exists \varepsilon(\lambda) \searrow 0$ so that

$$
\#\left\{\lambda_{j}:\left|\lambda_{j}-\lambda\right| \leq \varepsilon(\lambda)\right\}=O\left(\varepsilon(\lambda) \lambda^{n-1}\right),
$$

which implies the above nonclustering condition for $L^{p}$-resolvent bounds.
Specifically:

- Manifolds of nonpositive curvature (Bérard '78): $\varepsilon(\lambda)=1 / \log \lambda$
- Standard $n$-torus, $\mathbb{T}^{n}$ (Hlawka '50): $\varepsilon(\lambda)=\lambda^{-\sigma_{n}}$, $\sigma_{n}=-1+\frac{2}{n+1}$
- Duistermaat-Guillemin (\&lvrii) ('75) $\varepsilon(\lambda)=o(1)$ if $(M, g)$ has zero measure of periodic geodesics (is a generic condition).


## BSSY: Improvements of DKS-S in first two cases

Unlike the situation for the $n$-sphere, we can improve the earlier estimates of DKSS somewhat for i) manifolds of nonpositive curvature, and a bit more for ii) $\mathbb{T}^{n}$ :


Figure: Various regions $\mathcal{R}(M, g)$ for $\left(\Delta_{g}+\zeta\right)^{-1}$

Improved bounds for numerator in Sommerfeld-Green fnct

## Theorem (BSSY):

$$
\begin{equation*}
\|u\|_{L^{\frac{2 n}{n-2}}(M)} \leq C\left\|\left(\Delta_{g}+(\lambda+i \varepsilon(\lambda))^{2}\right) u\right\|_{L^{\frac{2 n}{n+2}}(M)}, \tag{5}
\end{equation*}
$$

if and only if

$$
\begin{equation*}
\left\|\sum_{\left|\lambda-\lambda_{j}\right| \leq \varepsilon(\lambda)} E_{j} f\right\|_{L^{\frac{2 n}{n-2}}(M)} \leq C \lambda \varepsilon(\lambda)\|f\|_{L^{\frac{2 n}{n+2}}(M)} . \tag{6}
\end{equation*}
$$

- $E_{j} f=<f, e_{j}>e_{j}(x)$ projection of $f$ onto $j$-th eigenspace. So above operator is proj of $f$ onto $\varepsilon(\lambda)$-spectral band about $\lambda$.
- CS ('88) earlier proved this estimate with $\varepsilon(\lambda)=1$ (sharp for sphere).
- Natural variant of (4), ie., Euclidean $L^{r} \rightarrow L^{s}, s^{\prime}=r$ bounds for

$$
\mathbb{1}_{[\lambda, \lambda+\varepsilon(\lambda)]]}\left(\sqrt{-\Delta_{\mathbb{R}^{n}}}\right) g=\int_{\left\{\xi \in \mathbb{R}^{n}:||\xi|-\lambda| \leq \varepsilon(\lambda)\right\}} e^{i \times \cdot \xi} \hat{g}(\xi) d \xi
$$

## Spectral projection bds $\Longrightarrow$ resolvent bds

Use the following variant of F.t. of Poisson summation kernel: If $P=\sqrt{-\Delta_{g}}, \zeta=(\lambda+i \mu)^{2}$,

$$
\begin{aligned}
\left(\Delta_{g}+(\lambda+i \mu)^{2}\right)^{-1} & =\frac{\operatorname{sgn} \mu}{i(\lambda+i \mu)} \int_{0}^{\infty} e^{i(\operatorname{sgn} \mu) \lambda t} e^{-|\mu| t}(\cos t P) d t \\
& =\ldots \int_{0}^{\infty} \beta(t) \ldots d t+\ldots \int_{0}^{\infty}(1-\beta(t)) \ldots d t
\end{aligned}
$$

with $C_{0}^{\infty} \ni \beta=1$ near origin.

- If $\beta$ small support, 1st term uniformly bounded for all $\zeta$ (local piece). Use Hadamard parametrix and Stein's osc int theorem.
- 2nd term is multiplier operator $m_{\lambda, \mu}(P)$ with

$$
m_{\lambda, \mu}(\tau)=O\left((1+|\lambda-\tau|)^{-N}\right)+O\left(|\mu|^{-1}\left(1+|\mu|^{-1}|\lambda-\tau|\right)^{-N}\right)
$$

and so can use old unit-band and improved $\mu=\varepsilon(\lambda)$-band bonds (and reverse $T T^{*}$ argument) to handle it.

## Improved spectral projection bounds of BSSY

Have

$$
\left\|\sum_{\left|\lambda-\lambda_{j}\right| \leq \varepsilon(\lambda)} E_{j} f\right\|_{L^{\frac{2 n}{n-2}}(M)} \leq C \lambda \varepsilon(\lambda)\|f\|_{L^{\frac{2 n}{n+2}(M)}} .
$$

with

- $\varepsilon(\lambda)=1 / \log \lambda$ if $(M, g)$ has nonpositive curvature.
- For $\mathbb{T}^{n}$ if $\varepsilon(\lambda)=\lambda^{-\sigma_{n}}$ with

$$
\sigma_{n}=\left\{\begin{array}{l}
\frac{85}{252} \approx 0.337, n=3 \\
\frac{2(n-1)}{n^{2}+2 n-2}, n \geq 4, \text { even } \\
\frac{2(n-1)}{n(n+1)}, n \geq 5, \text { odd }
\end{array}\right.
$$

- Tools: For $\mathbb{T}^{n}$ use Poisson summation formula and recent harmonic techniques developed by Bennett-Carbery-Tao, Bourgain-Guth...
For nonpos curv use Cartan-Hadamard thm to lift calculation up to universal cover \& variant of Poisson summation


## Results w/Shanglin Huang for hyperbolic space

$\mathbb{R}^{n} \mathrm{w} /$ metric of const curv $-\kappa, \kappa>0$ : In geod polar coords

$$
\begin{gathered}
\Delta_{-\kappa}=\partial_{r}^{2}+(n-1) \sqrt{\kappa} \operatorname{coth}(\sqrt{\kappa} r) \partial_{r}+\left(\sqrt{\kappa} \operatorname{csch}(\sqrt{\kappa} r)^{2} \Delta_{S^{n-1}}\right. \\
d V_{-\kappa}=\left(\frac{\sinh \sqrt{\kappa} r}{\sqrt{\kappa}}\right)^{n-1} d r d \theta, \quad 0<r<\infty,-\kappa<0
\end{gathered}
$$

Theorem
Let $n \geq 3$. Then for every $r, s$ as above $\exists C_{r, s}$ so that $(-\kappa=-1)$ if $u \in C_{0}^{\infty}$
$\|u\|_{L^{s}\left(\mathbb{H}^{n}, d V_{\mathbb{H}^{n}}\right)} \leq C_{r, s}\left\|\left(\left(\Delta_{\mathbb{H}^{n}}+\left(\frac{n-1}{2}\right)^{2}\right)+\zeta\right) u\right\|_{L^{r}\left(\mathbb{H}^{n}, d V_{\mathbb{H}^{n}}\right)},|\zeta| \geq 1$.
Also, have uniform bounds (indep of $-\kappa<0$ )
$\|u\|_{L^{s}\left(\mathbb{R}^{n}, d V_{-k}\right)} \leq C_{r, s}\left\|\left(\left(\Delta_{-\kappa}+\kappa\left(\frac{n-1}{2}\right)^{2}\right)+\zeta\right) u\right\|_{L^{r}\left(\mathbb{R}^{n}, d V_{-\kappa}\right)},|\zeta| \geq \kappa$.

Letting $\kappa \rightarrow 0_{+}$recover Euclidean estimates of KRS.

## Proof of resolvent estimates for constant neg curv

Use that for the shifted Laplacians $-\Delta_{-\kappa}-\kappa\left(\frac{n-1}{2}\right)^{2}$ (having spectrum $[0, \infty)$ ), have the following variants of Euclidean estimates (4) for $\sqrt{-\Delta_{-\kappa}-\kappa\left(\frac{n-1}{2}\right)^{2}}$ :

$$
\begin{equation*}
\left\|\mathbb{1}_{[\lambda, \lambda+\varepsilon]}\left(P_{\kappa}\right)\right\|_{L^{r}\left(d V_{-\kappa} \rightarrow L^{s}\left(d V_{-\kappa}\right)\right.} \leq C \varepsilon \lambda, \quad \lambda \geq \kappa . \tag{7}
\end{equation*}
$$

The condition that the spectral parameter be $\geq \kappa$ not necessary for $n=3$.
You prove this using Stein's analytic interpolation argument used to prove Stein-Tomas and explicit formulae for functions of $P_{\kappa}$ (coming from explicit formulae for fund solution of wave operators)

Use (7) to deal with large frequencies $\lambda$ in spectral decomposition of resolvent operators

For low freq: Very handy to use Sobolev estimates for unshifted Laplacian: $\left(-\Delta_{\mathbb{H}^{n}}\right)^{\alpha / 2}: L^{p}(\alpha)\left(\mathbb{H}^{n}\right) \rightarrow L^{q(\alpha)}\left(\mathbb{H}^{n}\right)$ (exponents as in $\mathbb{R}^{n}$ ) from Cowling-Giulini-Meda ('93)

## Constant positive curvature: round spheres

Also can prove continuous family of uniform estimates for spheres of constant curvature $\kappa>0$ that imply the KRS estimates.


Do this by proving uniform $L^{r} \rightarrow L^{s}$ bounds for proj onto spherical harmonics of degree $k$. Recover $\mathbb{R}^{n}$ bounds by letting $\kappa \rightarrow 0_{+}$ Need $\left\|H_{k}\right\|_{L^{r}\left(S^{n}\right) \rightarrow L^{s}\left(S^{n}\right)}$ (proj onto spherical harmonics of deg $k$.)

For this, need to strengthen classical Darboux formula for asymptotics of projection kernel (i.e., "zonal functions" on $S^{n}$ )
Use periodicity of wave group $t \rightarrow e^{i t \sqrt{-\Delta_{S^{n}}+\left(\frac{n-1}{2}\right)^{2}}}$.

