

$$R^2 g = 2 du dR + \frac{dS d\bar{T}}{(1 + |S|^2)^2} + R^2 du^2.$$

S matrix.

Gravitational phase space:

$$R^2 ds^2 = du dR + \frac{dS d\bar{T}}{(1 + |S|^2)^2}.$$

$$+ \frac{R \sigma dt^2 + c.c.}{(1 + |S|^2)^2} \rightarrow O(R)$$



Asymp. sym. data $\sigma(u, S, \bar{S})$ "shear".

H^+ = phase space for σ , and has a symplectic structure.

Symmetries:

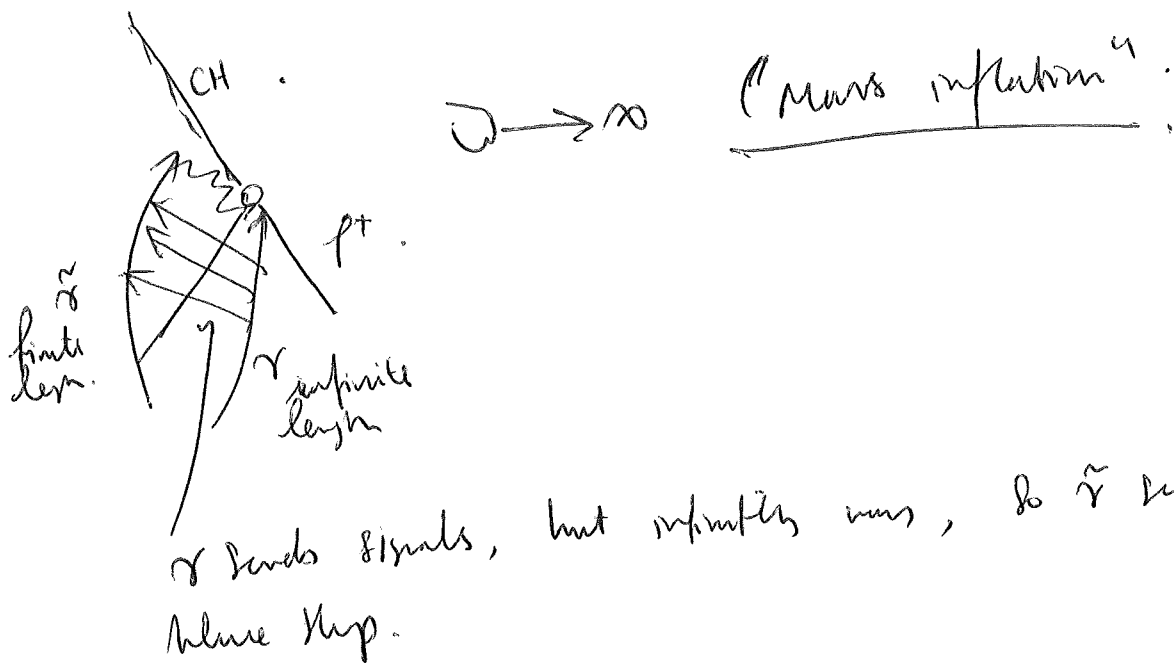
$\sigma_- \in \mathcal{H}_-$ solves Gravitational problem

$S[\sigma_-]$ = renormalized Einstein action.

Scattering $\sigma_+ = \mathcal{I}(\sigma_-)$

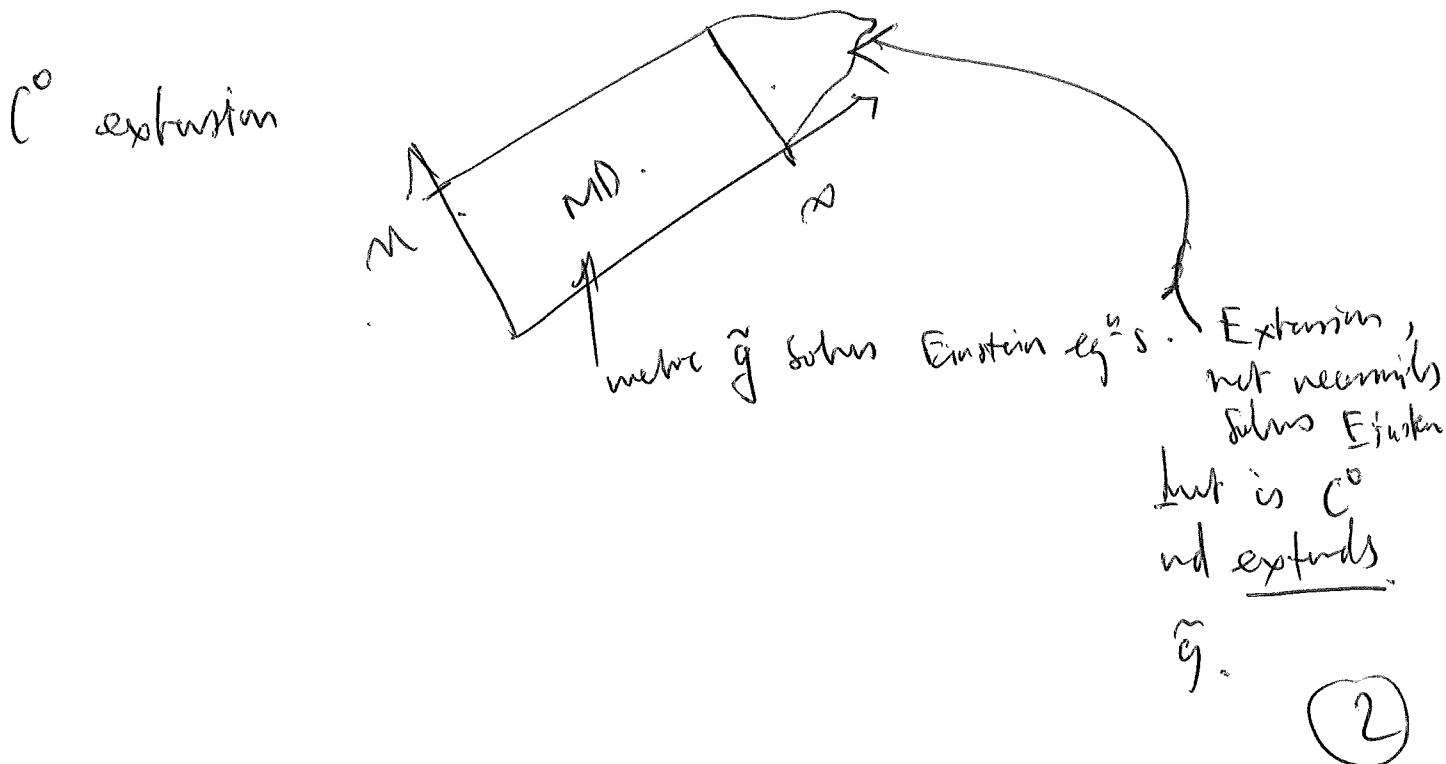
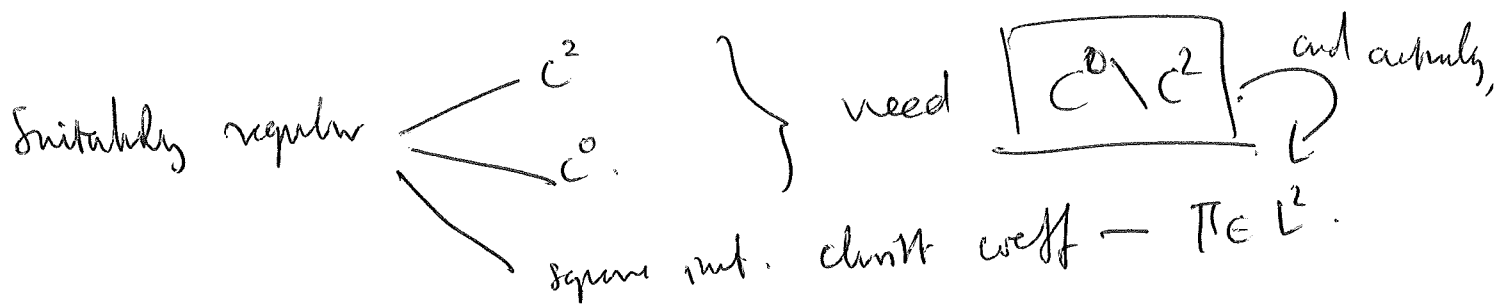
$$\Omega(\sigma_+, \sigma'_-) = \left. \frac{-d}{d\varepsilon} S[\sigma_- + \varepsilon \sigma'_-] \right|_{\varepsilon=0}.$$

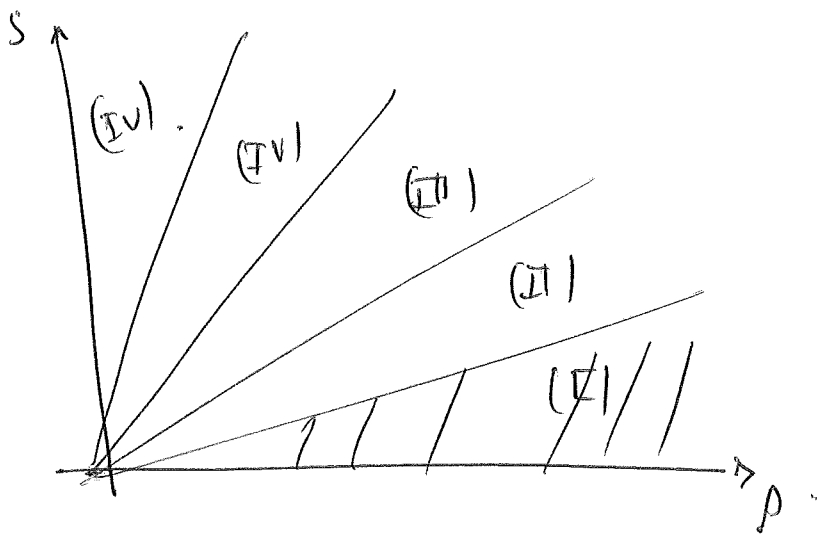
69; Argument by Penrose & ...? $\Lambda = 0$.



Strong Cosmic Censorship: as a "suitably regular" C^0 extension.

Generically, MD is inextendible! horizon ruled.

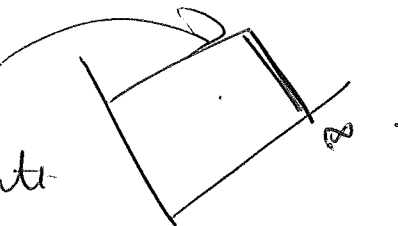




$$\partial_{\mu} \varphi \sim \mu^s$$

$$1 < \beta = \frac{h_-}{h_+} = \frac{\text{surface} \rightarrow \infty}{\text{quantity} \rightarrow h^3}$$

(I) $\omega \rightarrow \infty, \nu \rightarrow \infty$



$\hat{N} = f(\nu)$ constant
 number.

$$R_{\text{min}}^2 \gtrsim \omega^2 + O(\beta) \rightarrow \infty \Rightarrow \nexists c^2 \text{ ext.}$$

\Rightarrow get $\Pi \notin L^2_{\text{loc}}$.

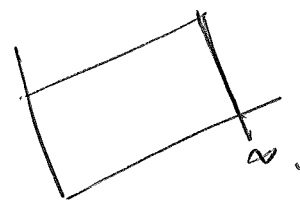
(II). $\omega \rightarrow \infty$ in $|\frac{\partial \varphi}{\partial R}| \rightarrow \infty$, $R_{\text{min}} \gtrsim \left(\frac{\partial \varphi}{\partial R}\right)^2$.

(III) - (IV) no mass inflation $\Leftrightarrow |\omega| < c \Rightarrow \Pi \in L^2_{\text{loc}}$.

In (IV) main thing to show $|\frac{\partial \varphi}{\partial R}| \leq C$.

\Rightarrow go all the way to ∞ .

with $|\frac{\partial \varphi}{\partial R}| \leq C$.



\exists classical solⁿ of EMS in $C^1 \setminus C^2$, $R_{\text{min}} \in C^0$.